

Detection and Imaging of Multiple Ground Moving Targets using Ultra-Narrowband Continuous-Wave SAR

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ABSTRACT

We present a novel method for ground moving target detection and imaging using a SAR system transmitting ultra-narrowband continuous waveforms. We develop a new forward model that relates the velocity as well as reflectivity information at each location to a correlated received signal. We reconstruct moving target images by a filtered-backprojection method. We use the image contrast as a metric to detect moving targets and to determine their velocities. The method results in well-focused reflectivity images of moving targets and their velocity estimates regardless of the target location, speed, and velocity direction. We present numerical experiments to verify our method.

Keywords: Moving Target, Detection, Imaging, Synthetic Aperture, Continuous Wave (CW), Doppler

1. INTRODUCTION

Ground moving target detection and imaging using synthetic aperture radar (SAR) has received considerable attention in the past two decades.¹⁻¹⁷ However, all the existing techniques are for traditional synthetic aperture radar (SAR) systems, which accomplish high-range resolution imaging of static scenes and moving targets by transmitting wideband waveforms. We consider the problem of synthetic aperture moving target detection and imaging using ultra-narrowband CW waveforms. Such a SAR system requires relatively simple and low-cost transmitter, and in some cases does not need a dedicated transmitter. Existing radio frequency signals, such as radio, television signals, WiFi signals, etc. can be used as the transmission sources. Additionally, ultra-narrowband continuous waveforms have high Doppler resolution and are capable of capturing the velocity information of the moving targets as compared to the high range resolution waveforms used by the traditional SAR systems.

In [18], we presented a novel synthetic aperture imaging method of stationary scenes that takes advantage of the high Doppler resolution of the transmitted ultra-narrowband continuous waveforms. In this paper, we consider the moving target detection and imaging problem using such SAR systems. Our approach exploits the high Doppler resolution of the transmitted waveforms to reconstruct the reflectivity (position), as well as to estimate the velocity field of moving targets. We develop a novel model that relates the velocity field and the reflectivity of the scene to a correlated received signal and a corresponding FBP-type novel image formation method. A set of FBP images are reconstructed with hypothesized velocities that cover the expected velocity range. Moving targets are well focused in the reflectivity images corresponding to their velocities. We use the image contrast as a metric to detect targets and to determine their velocities. We present numerical experiments to demonstrate the performance of the moving target detection and imaging method. To the best of our knowledge, our method is the first in the literature that addresses the ground moving target detection and imaging using such SAR systems.

The organization of the paper is as follows. In Section 2, we present the model for the received signal and develop the forward model for moving target imaging using ultranarrow-band CW waveforms. In Section 3,

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we develop an FBP-type image formation method to reconstruct the reflectivity of the scene and a contrast-maximization based velocity estimation method. In Section 4, we present numerical simulations to demonstrate the performance of the proposed moving target imaging method. Section 5 concludes our paper.

We use the following notational conventions throughout the paper. The bold Roman, bold italic and Roman lower-case letters are used to denote variables in \mathbb{R}^3 , \mathbb{R}^2 and \mathbb{R} , respectively, i.e., $\mathbf{z} = (z, z) \in \mathbb{R}^3$, with $z \in \mathbb{R}^2$ and $z \in \mathbb{R}$. The calligraphic letters (\mathcal{F}, \mathcal{K} etc.) are used to denote operators.

2. FORWARD MODEL FOR MOVING TARGET IMAGING

2.1. Received Signal

Let $\mathbf{x} = (\mathbf{x}, \psi(\mathbf{x})) \in \mathbb{R}^3$ denote the earth's surface, where $\mathbf{x} \in \mathbb{R}^2$ and $\psi : \mathbb{R}^2 \rightarrow \mathbb{R}$ is a known function for the ground topography. Let $\mathbf{v}_\mathbf{x}$ denote the velocity of the moving scatterer located at \mathbf{x} at $t = 0$. Then,

$$\mathbf{v}_\mathbf{x} = [\mathbf{v}_\mathbf{x}, \nabla_\mathbf{x}\psi(\mathbf{x}) \cdot \mathbf{v}_\mathbf{x}] \quad (1)$$

where $\mathbf{v}_\mathbf{x} \in \mathbb{R}^2$ is the two dimensional velocity and $\nabla_\mathbf{x}\psi(\mathbf{x}) = [\frac{\partial\psi}{\partial x_1} \quad \frac{\partial\psi}{\partial x_2}]$. Note that $\mathbf{v}_\mathbf{x}$ may be t dependent. In the rest of our discussion, we make the assumption that the scatterers are moving linearly and therefore the velocity $\mathbf{v}_\mathbf{x}$ is time independent. Furthermore, we assume that the functional form of $\mathbf{v}_\mathbf{x}$ does not depend on \mathbf{x} . It is simply a constant velocity for each \mathbf{x} . Thus, at time t , the scatter is located at $\mathbf{z} = \mathbf{x} + \mathbf{v}_\mathbf{x}t$.

We assume that the scattering takes place in a thin region near the ground. Thus, the reflectivity function of the scene at time $t = 0$ has the form

$$V(\mathbf{x}) = \rho(\mathbf{x})\delta(\mathbf{x} - \psi(\mathbf{x})). \quad (2)$$

At time t , taking into account the movement of the scatters, the reflectivity function of the moving scatterers translates as

$$V(\mathbf{z} - \mathbf{v}_\mathbf{x}t) \simeq \rho(\mathbf{z} - \mathbf{v}_\mathbf{x}t)\delta(z_3 - \psi(\mathbf{z}))\delta(v_3 - \nabla_\mathbf{z}\psi(\mathbf{z}) \cdot \mathbf{v}_\mathbf{x}). \quad (3)$$

Let $\gamma_T(t)$ and $\gamma_R(t)$ be the transmitter and receiver trajectories, respectively; and $p(t)$ denotes the transmitted waveform and $r(t)$ denote the received signal. We consider a finite and relatively short receiving time-window starting at time $t = s$.

Using the scalar wave equation along with (3) and under the Born approximation, taking into account the fact that the velocities of the antennas and targets are much less than the speed of light; and under the assumption that the range variation due to the movement of the antennas and the scene is much less than the the transmitter-to-scene and scene-to-receiver range, we have¹⁸

$$r(t + s) = -\omega_0^2 \int \frac{\tilde{p}(\alpha t - \tau + s)e^{i\omega_0(\alpha t - \tau + s)}q(\mathbf{x}, \mathbf{v})}{(4\pi)^2 G_{TR}(s, \mathbf{x}, \mathbf{v})} d\mathbf{x}d\mathbf{v} \quad (4)$$

where $q(\mathbf{x}, \mathbf{v}) = \rho(\mathbf{x})\delta(\mathbf{v} - \mathbf{v}_\mathbf{x})$ is the *phase-space reflectivity function* of the moving scene. Note that a narrowband waveform is assumed in (4), i.e., $p(t) = e^{i\omega_0 t}\tilde{p}(t)$ where ω_0 denotes the carrier frequency and $\tilde{p}(t)$ is the complex envelope of p , which is slow varying as a function of t as compared to $e^{i\omega_0 t}$.

The time dilation α in (4) is given by

$$\alpha = \frac{1 - \gamma_R(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} \cdot \dot{\gamma}_R(s)/c_0}{1 + \gamma_T(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} \cdot \dot{\gamma}_T(s)/c_0} \cdot \frac{1 + \gamma_T(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} \cdot \mathbf{v}/c_0}{1 - \gamma_R(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} \cdot \mathbf{v}/c_0}, \quad (5)$$

the time delay τ is given by

$$\tau \approx [|\gamma_T(s) - \widehat{(\mathbf{x} + \mathbf{v}s)}| + |\gamma_R(s) - \widehat{(\mathbf{x} + \mathbf{v}s)}|]/c_0 - [(\gamma_T(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} + \gamma_R(s) - \widehat{(\mathbf{x} + \mathbf{v}s)}) \cdot \mathbf{v}s]/c_0 \quad (6)$$

and G_{TR} is the product of the geometrical spreading factors given by

$$G_{TR}(s, \mathbf{x}, \mathbf{v}) = |\boldsymbol{\gamma}_R(s) - (\mathbf{x} + \mathbf{v}s)| |\mathbf{x} + \mathbf{v}s - \boldsymbol{\gamma}_T(s)|. \quad (7)$$

Since \tilde{p} is a slow-varying function of time, we approximate $\tilde{p}(\alpha t) \approx \tilde{p}(t)$ in the rest of our discussion.

Note that the time dilation factor α can be approximated as $\alpha = 1 + \beta$ where $f_0\beta$ ($f_0 = \omega_0/2\pi$) represents the total Doppler frequency induced by the relative radial motion of the antennas and the moving scatters. We refer to it as the *bistatic Doppler frequency for moving targets* and denote it with $f_d(s, \mathbf{x}, \mathbf{v})$, i.e.,

$$f_d(s, \mathbf{x}, \mathbf{v}) = \frac{f_0}{c_0} [\boldsymbol{\gamma}_T(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} \cdot (\dot{\boldsymbol{\gamma}}_T(s) - \mathbf{v}) + \boldsymbol{\gamma}_R(s) - \widehat{(\mathbf{x} + \mathbf{v}s)} \cdot (\dot{\boldsymbol{\gamma}}_R(s) - \mathbf{v})]. \quad (8)$$

Note that in (7) and (8), $\mathbf{x} = [\mathbf{x}, \boldsymbol{\psi}(\mathbf{x})]$ and $\mathbf{v} = [\mathbf{v}, \nabla_{\mathbf{x}}\boldsymbol{\psi}(\mathbf{x}) \cdot \mathbf{v}]$.

2.2. Forward Model

We define the correlation of the received signal given in (4) with a scaled or frequency-shifted version of the transmitted signal over a finite time window as follows:

$$d(s, \mu) = \int r(t+s)p^*(\mu t)\phi(t)dt \quad (9)$$

for some $s \in \mathbb{R}$ and $\mu \in \mathbb{R}^+$, where $\phi(t)$, $t \in [0, T_\phi]$ is a smooth windowing function with a finite support.

Substituting (4) into (9), we obtain

$$d(s, \mu) = \int \frac{e^{i\omega_0(\alpha-\mu)t} e^{i\omega_0(s-\tau)}}{(4\pi)^2 G_{TR}(s, \mathbf{x}, \mathbf{v})} \omega_0^4 \tilde{p}(t-\tau+s) \tilde{p}^*(t) q(\mathbf{x}, \mathbf{v}) d\mathbf{x} d\mathbf{v} dt. \quad (10)$$

Note that $\tilde{p}(\mu t) \approx \tilde{p}(t)$ is used in (10).

We define the forward modeling operator, \mathcal{F} , as follows:

$$\begin{aligned} d(s, \mu) &\approx \mathcal{F}[q](s, \mu) \\ &:= \int e^{-i\phi(t, \mathbf{x}, \mathbf{v}, s, \mu)} A(t, \mathbf{x}, \mathbf{v}, s, \mu) q(\mathbf{x}, \mathbf{v}) d\mathbf{x} d\mathbf{v} dt \end{aligned} \quad (11)$$

where

$$\phi(t, \mathbf{x}, \mathbf{v}, s, \mu) = 2\pi f_0 t [(\mu - 1) + f_d(s, \mathbf{x}, \mathbf{v})/f_0] \quad (12)$$

$$A(t, \mathbf{x}, \mathbf{v}, s, \mu) = \frac{\tilde{p}(t-\tau+s) \tilde{p}^*(t) e^{i\omega_0(s-\tau)} \omega_0^4}{(4\pi)^2 G_{TR}(s, \mathbf{x}, \mathbf{v})} \quad (13)$$

We assume that for some m_A , A satisfies the inequality

$$\sup_{(t, \mu, s, \mathbf{x}, \mathbf{v}) \in \mathcal{U}} \left| \partial_t^{\alpha_t} \partial_\mu^{\alpha_\mu} \partial_s^{\beta_s} \partial_{x_1}^{\epsilon_1} \partial_{x_2}^{\epsilon_2} \partial_{v_1}^{\epsilon_1} \partial_{v_2}^{\epsilon_2} A(t, \mathbf{x}, \mathbf{v}, s, \mu) \right| \leq C_A (1 + t^2)^{(m_A - |\alpha_t|)/2} \quad (14)$$

where \mathcal{U} is any compact subset of $\mathbb{R} \times \mathbb{R}^+ \times \mathbb{R} \times \mathbb{R}^2 \times \mathbb{R}^2$, and the constant C_A depends on \mathcal{U} , $\alpha_t, \mu, \beta_s, \epsilon_{1,2}, \epsilon_{1,2}$. This assumption is needed in order to make various stationary phase calculations hold.

3. IMAGE FORMATION

We assume that the velocity is constant, say \mathbf{v}_h , and reconstruct a set of two-dimensional reflectivity images in position space for a range of hypothesized velocities $\mathbf{v}_h \in \Lambda_v$. We refer to this image as the \mathbf{v}_h -reflectivity image and form it by an FBP-type imaging operator. We design the filter to ensure that the reconstructed reflectivity for the scatterer at \mathbf{x} has the correct strength whenever the hypothesized velocity is equal to the true target velocity, i.e., $\mathbf{v}_h = \mathbf{v}_x$. From this set of images, we estimate the velocity of the targets using a figure of merit that measures the degree to which the images are focused. The reflectivity images corresponding to the estimated velocities provide focused images of the moving targets present in the scene.

3.1. FBP- operator

We form the \mathbf{v}_h -reflectivity image $q_{\mathbf{v}_h}(\mathbf{z})$ for a fixed hypothesized velocity $\mathbf{v}_h = [\mathbf{v}_h, \nabla_{\mathbf{z}}\psi(\mathbf{z}) \cdot \mathbf{v}_h]$ by filtering and backprojecting the data onto the position-space iso-Doppler contour $F_{\mathbf{v}_h}(s, \mu)$:

$$\begin{aligned} q_{\mathbf{v}_h}(\mathbf{z}) &:= \mathcal{K}_{\mathbf{v}_h}[d](\mathbf{z}) \\ &= \int e^{i\phi(t, \mathbf{z}, \mathbf{v}_h, s, \mu)} Q_{\mathbf{v}_h}(\mathbf{z}, t, s) d(s, \mu) dt ds d\mu, \end{aligned} \quad (15)$$

where $\mathcal{K}_{\mathbf{v}_h}$ is the filtered-backprojection operator for the fixed velocity \mathbf{v}_h and $Q_{\mathbf{v}_h}$ is the filter to be determined below. Note that \mathbf{v}_h is a fixed parameter for $Q_{\mathbf{v}_h}$.

We assume that for some $m_{Q_{\mathbf{v}_h}}$, $Q_{\mathbf{v}_h}$ satisfies the inequality

$$\sup_{(t, \mu, s, \mathbf{z}, \mathbf{v}_h) \in \mathcal{U}} \left| \partial_t^{\alpha_t} \partial_\mu^{\alpha_\mu} \partial_s^{\beta_s} \partial_{z_1}^{\epsilon_1} \partial_{z_2}^{\epsilon_2} Q_{\mathbf{v}_h}(t, \mathbf{z}, \mathbf{v}_h, s, \mu) \right| \leq C_{Q_{\mathbf{v}_h}} (1 + t^2)^{(m_{Q_{\mathbf{v}_h}} - |\alpha_t|)/2} \quad (16)$$

where \mathcal{U} is any compact subset of $\mathbb{R} \times \mathbb{R}^+ \times \mathbb{R} \times \mathbb{R}^2$, and the constant $C_{Q_{\mathbf{v}_h}}$ depends on $\mathcal{U}, \alpha_{t, \mu}, \beta_s, \epsilon_{1,2}$. Under the assumption (16), (15) defines $\mathcal{K}_{\mathbf{v}_h}$ as a Fourier integral operator.

The kernel of the image fidelity operator $\mathcal{K}_{\mathbf{v}_h} \mathcal{F}$ is the *Point Spread Function* (PSF) of the two-dimensional reflectivity imaging operator for the hypothesized velocity \mathbf{v}_h with respect to the true velocity \mathbf{v}_x , which is given by

$$L_{\mathbf{v}_h}^{\mathbf{v}_x}(\mathbf{z}, \mathbf{x}) = \int e^{i[\phi(t, \mathbf{z}, \mathbf{v}_h, s, \mu) - \phi(t, \mathbf{x}, \mathbf{v}_x, s, \mu)]} Q_{\mathbf{v}_h}(\mathbf{z}, t, s) A(t', \mathbf{x}, \mathbf{v}_x, s, \mu) dt ds d\mu dt'. \quad (17)$$

Applying the stationary phase theorem to approximate the t' and μ integrations in (17) * and substituting the results back into (17), we get the kernel of the image fidelity operator $\mathcal{K}_{\mathbf{v}_h} \mathcal{F}$:

$$L_{\mathbf{v}_h}^{\mathbf{v}_x}(\mathbf{z}, \mathbf{x}) \approx \int e^{i2\pi t[f_d(s, \mathbf{z}, \mathbf{v}_h) - f_d(s, \mathbf{x}, \mathbf{v}_x)]} Q_{\mathbf{v}_h}(\mathbf{z}, t, s) A(t, \mathbf{x}, \mathbf{v}_x, s, 1 - f_d(s, \mathbf{x}, \mathbf{v}_x)/f_0) dt ds. \quad (18)$$

To simplify our notation, we let

$$A(t, \mathbf{x}, \mathbf{v}_x, s) = A(t, \mathbf{x}, \mathbf{v}_x, s, 1 - f_d(s, \mathbf{x}, \mathbf{v}_x)/f_0). \quad (19)$$

We determine $Q_{\mathbf{v}_h}$ so that the PSF of the two-dimensional reflectivity imaging operator, $L_{\mathbf{v}_h}^{\mathbf{v}_x}(\mathbf{z}, \mathbf{x})$ is as close as possible to the Dirac-delta function, $\delta(\mathbf{z} - \mathbf{x})$ for $\mathbf{v}_h = \mathbf{v}_x$, i.e., whenever the reflectivity at \mathbf{z} is reconstructed for the correct target velocity \mathbf{v}_x . We assume that at the correct target velocity, the flight trajectory and the illumination pattern are chosen such that the only contribution to $L_{\mathbf{v}_h}^{\mathbf{v}_x}(\mathbf{z}, \mathbf{x})$ comes from those points $\mathbf{z} = \mathbf{x}$.

Thus, we linearize $f_d(s, \mathbf{z}, \mathbf{v}_h)$ around $\mathbf{z} = \mathbf{x}$ for $\mathbf{v}_h = \mathbf{v}_x$ and approximate

$$f_d(s, \mathbf{z}, \mathbf{v}_h) - f_d(s, \mathbf{x}, \mathbf{v}_h) \approx (\mathbf{z} - \mathbf{x}) \cdot \nabla_{\mathbf{z}} f_d(s, \mathbf{z}, \mathbf{v}_h). \quad (20)$$

Thus, (18) becomes

$$L_{\mathbf{v}_h}^{\mathbf{v}_x}(\mathbf{z}, \mathbf{x}) = \int e^{it(\mathbf{z} - \mathbf{x}) \cdot \Xi_{\mathbf{v}_h}(s, \mathbf{z})} Q_{\mathbf{v}_h}(\mathbf{z}, t, s) A(t, \mathbf{z}, \mathbf{v}_h, s) dt ds \quad (21)$$

where

$$\Xi_{\mathbf{v}_h}(s, \mathbf{z}) = 2\pi \nabla_{\mathbf{z}} f_d(s, \mathbf{z}, \mathbf{v}_h). \quad (22)$$

For each \mathbf{z} , we make the following change of variables:

$$(t, s) \rightarrow \boldsymbol{\xi} = t \Xi_{\mathbf{v}_h}(s, \mathbf{z}) \quad (23)$$

*The determinant of the Hessian of Φ_k is $(2\pi)^2 f_0^2$. Thus, the stationary points are non-degenerate.

and write (21) as follows:

$$L_{\mathbf{v}_h}^{\mathbf{v}_h}(\mathbf{z}, \mathbf{x}) = \int_{\Omega_{\mathbf{v}_h, \mathbf{z}}} e^{i(\mathbf{z}-\mathbf{x}) \cdot \boldsymbol{\xi}} Q_{\mathbf{v}_h}(\mathbf{z}, \boldsymbol{\xi}) A(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi}) \eta(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi}) d\boldsymbol{\xi} \quad (24)$$

where

$$Q_{\mathbf{v}_h}(\mathbf{z}, \boldsymbol{\xi}) = Q_{\mathbf{v}_h}(\mathbf{z}, t(\boldsymbol{\xi}), s(\boldsymbol{\xi})), \quad (25)$$

$$A(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi}) = A(t(\boldsymbol{\xi}), \mathbf{z}, \mathbf{v}_h, s(\boldsymbol{\xi})) \quad (26)$$

and

$$\eta(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi}) = \frac{\partial(t, s)}{\partial \boldsymbol{\xi}} = |t|^{-1} \left| \det \begin{bmatrix} \boldsymbol{\Xi}_{\mathbf{v}_h}(s, \mathbf{z}) \\ \partial_s \boldsymbol{\Xi}_{\mathbf{v}_h}(s, \mathbf{z}) \end{bmatrix} \right|^{-1} \quad (27)$$

is the determinant of the Jacobian that comes from the change of variables given in (23).

The domain of integration in (24) is given by

$$\Omega_{\mathbf{v}_h, \mathbf{z}} = \{\boldsymbol{\xi} = t \boldsymbol{\Xi}_{\mathbf{v}_h}(s, \mathbf{z}) \mid A(t, \mathbf{z}, \mathbf{v}_h, s) \neq 0, \quad t, s \in \mathbb{R}\}. \quad (28)$$

We refer to $\Omega_{\mathbf{v}_h, \mathbf{z}}$ as the *data collection manifold* at \mathbf{z} for $\mathbf{v}_h = \mathbf{v}_x$. This set determines many of the properties of the reconstructed reflectivity image when $\mathbf{v}_h = \mathbf{v}_x$.

To approximate the point spread function $L_{\mathbf{v}_h}^{\mathbf{v}_h}(\mathbf{z}, \mathbf{x})$ in (24) with the Dirac-delta function, we choose the filter as follows:

$$Q_{\mathbf{v}_h}(\mathbf{z}, \boldsymbol{\xi}) = \frac{\chi_{\Omega_{\mathbf{v}_h, \mathbf{z}}}}{\eta(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi})} \frac{A^*(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi})}{|A(\mathbf{z}, \mathbf{v}_h, \boldsymbol{\xi})|^2} \quad (29)$$

where $\chi_{\Omega_{\mathbf{v}_h, \mathbf{z}}}$ is a smooth cut-off function that prevents division by zero in (29).

3.2. Determination of the moving target velocity

The filtered-backprojection (FBP) of $d(s, \mu)$ results in a set of reflectivity images $\tilde{q}_{\mathbf{v}_h}$ in the two-dimensional \mathbf{z} space for each velocity value \mathbf{v}_h for a range of velocities that is suitably chosen for ground moving targets. When the hypothesized velocity \mathbf{v}_h is equal to the correct velocity \mathbf{v}_z , the corresponding reflectivity image is expected to be well-focused. We measure the degree to which the reflectivity images are focused with the image contrast and generate a contrast-image defined as follows:

$$I(v_{h1}, v_{h2}) = \mathcal{C}[q_{\mathbf{v}_h}] \quad (30)$$

where \mathcal{C} denotes the contrast operator³ and $\mathbf{v}_h = (v_{h1}, v_{h2})$ is used as the index of the contrast image. We determine the velocity of the moving scatterers by localizing the local maxima in the contrast-image $I(\mathbf{v}_h)$.

4. NUMERICAL SIMULATIONS

We considered a scene of size $[0, 11] \times [0, 11]$ km² with flat topography centered at $[11, 11, 0]$ km. The scene was discretized into 128×128 pixels, where $[0, 0, 0]$ m and $[11, 11, 0]$ km correspond to the pixels (1, 1) and (128, 128), respectively. Fig. 1 shows the scene with a static extended target and multiple moving targets along with their corresponding velocities.

We assumed that the transmitter and receiver were traversing a circular trajectory given by $\gamma_C(s) = (11 + 11 \cos(s), 11 + 11 \sin(s), 6.5)$ km. Let $\gamma_T(s)$ and $\gamma_R(s)$ denote the trajectories of the two receivers. We set $\gamma_T(s) = \gamma_C(s)$ and $\gamma_R(s) = \gamma_1(s - \frac{\pi}{4})$. Note that the variable s in γ_C is equal to $\frac{V}{R}t$ where V is the speed of the receiver, and R is the radius of the circular trajectory. We set the speed of the two receivers to 261 m/s.

We assumed that the transmitter transmitted a single-frequency continuous waveform operating at $f_0 = \omega_0/2\pi = 800$ MHz. We used (4) and (9) to generate the data. The length of the signal was set to 0.1707 s. The circular trajectory was uniformly sampled into 2048 points.

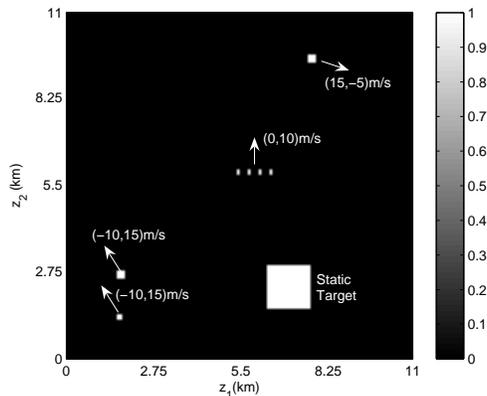


Figure 1: Scene considered in the numerical simulations.

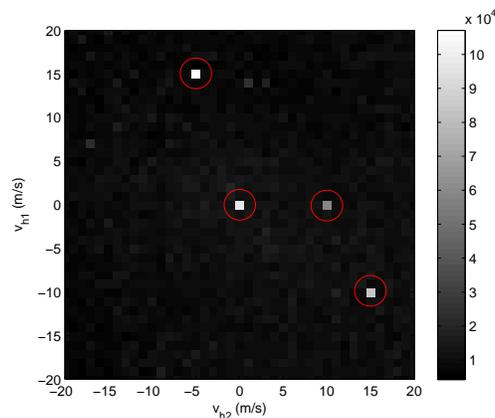


Figure 2: The contrast-image formed by the contrast of the images reconstructed with each hypothesized velocity.

We assumed that the velocity of the targets is in the range of $[-20, 20] \times [-20, 20]$ m/s and reconstructed $q_{v_h}(z)$ images via the FBP and image-contrast optimization method as described in Section 3. We see from Fig. 2 that there are four dominant peaks marked with red circles in the contrast-image. This indicates that there are four different velocities associated with the targets in the scene. The velocities where the peaks are located are equal to the velocities of the targets we assumed in the simulations.

Fig. 3 presents the reconstructed images using the four estimated velocities. We observe that the targets are well-focused in the image formed using the correct velocity associated with each target. Note that Fig. 3(a) is the image reconstructed with $\mathbf{v}_h = [0, 0, 0]$ m/s. In this case, the moving target imaging method described here is equivalent to the static target imaging method in.¹⁹ As expected only the static target is well-focused in Fig. 3(a).

5. CONCLUSION

In this paper, we considered the problem of synthetic aperture radar imaging of moving targets using ultra-narrowband transmitted waveforms. We presented a received signal model for a dynamic scene. We developed a novel forward model for image formation. We developed an associated novel FBP and image-contrast optimization based image formation method to estimate the velocities of multiple targets and their corresponding reflectivity images. We presented numerical simulations to verify the theoretical results. A complete analysis of the performance of our imaging method and demonstration of its performance in more realistic scenarios will be the focus of our future work.

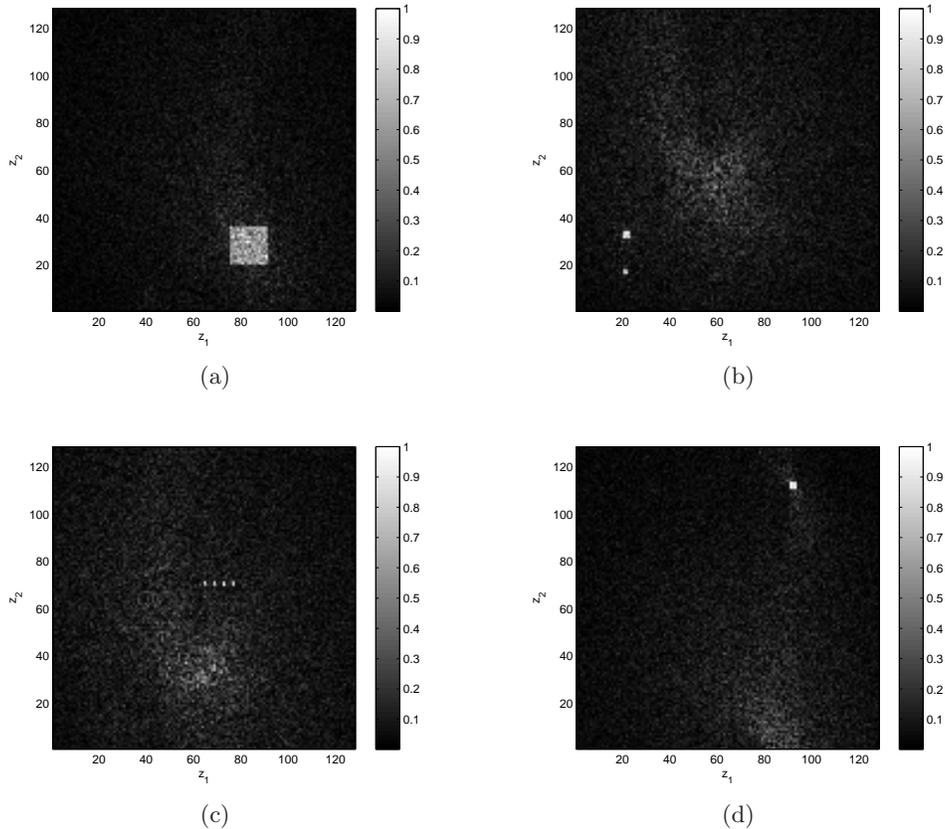


Figure 3: Reconstructed images with the target velocities: (a) $\mathbf{v}_h = [0, 0, 0]$ m/s; (b) $\mathbf{v}_h = [-10, 15, 0]$ m/s; (c) $\mathbf{v}_h = [0, 10, 0]$ m/s; (d) $\mathbf{v}_h = [15, -5, 0]$ m/s.

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