

# Russel Ohl – Inventor of a p-n junction

The most important professor's function  
Is to explain the workings of a p-n **Junction**

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From <http://nina.ecse.rpi.edu/shur/Alphab.htm>

1940

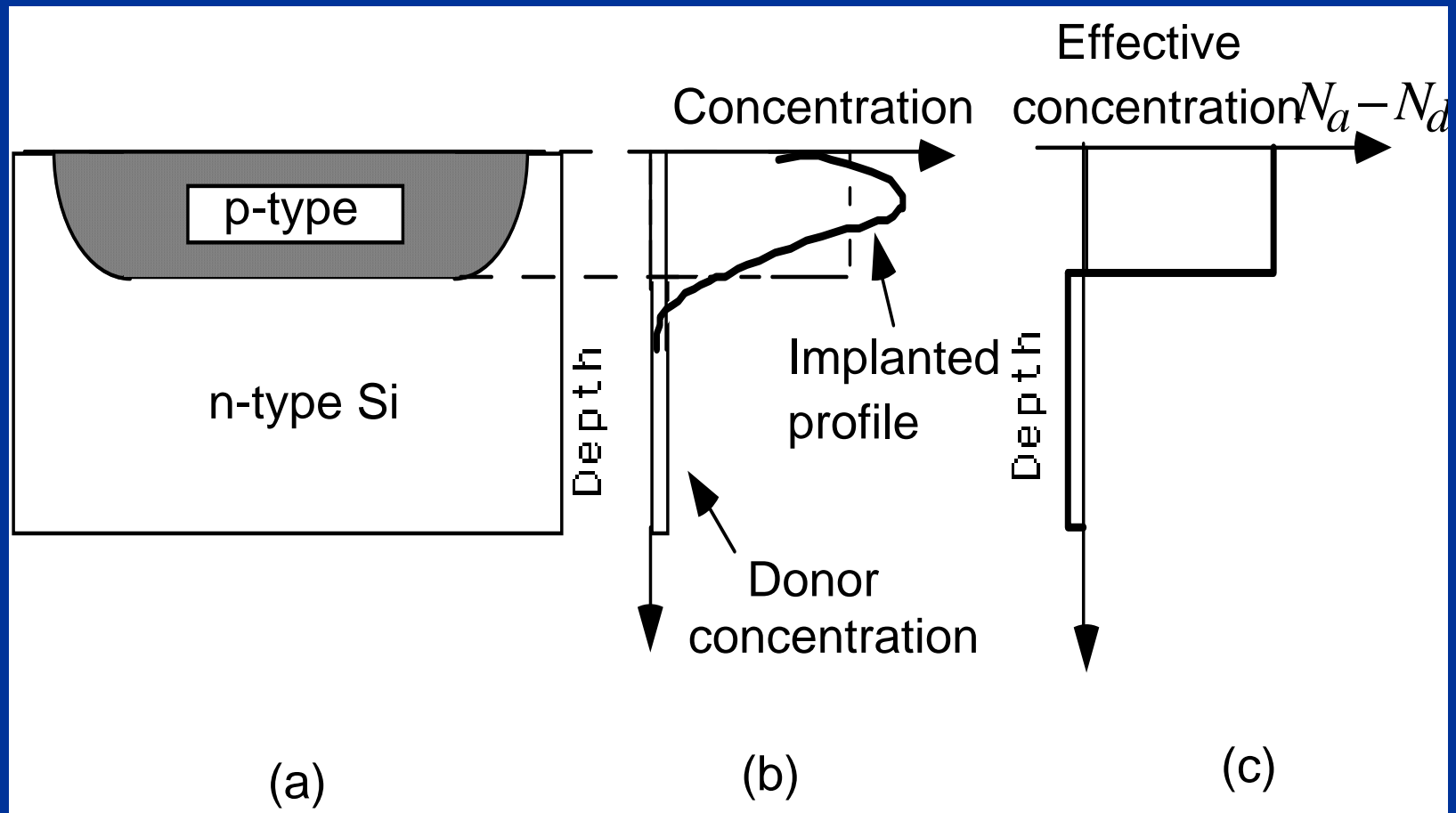


# Professor Karmalkar ITT Madras

<http://www.youtube.com/watch?v=IMoJUqDISQs>

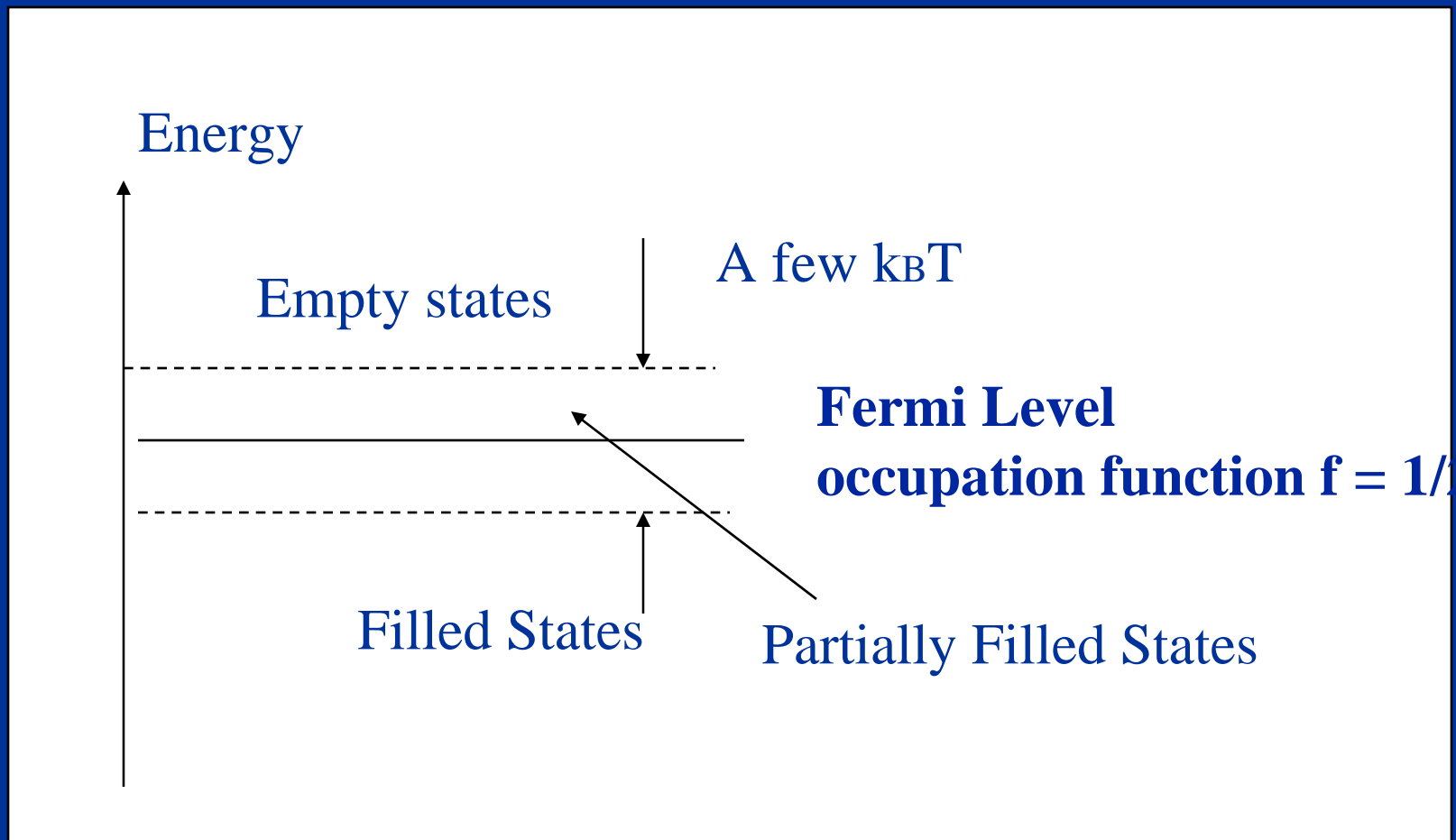


# p-n Junction



**Boron implantation into n-type Si  
to produce a p region**

# Fermi Level



## Fermi Level in Equilibrium

In equilibrium, the electron and hole current densities must be zero, and the quasi-Fermi levels of electrons and holes coincide with the position of the

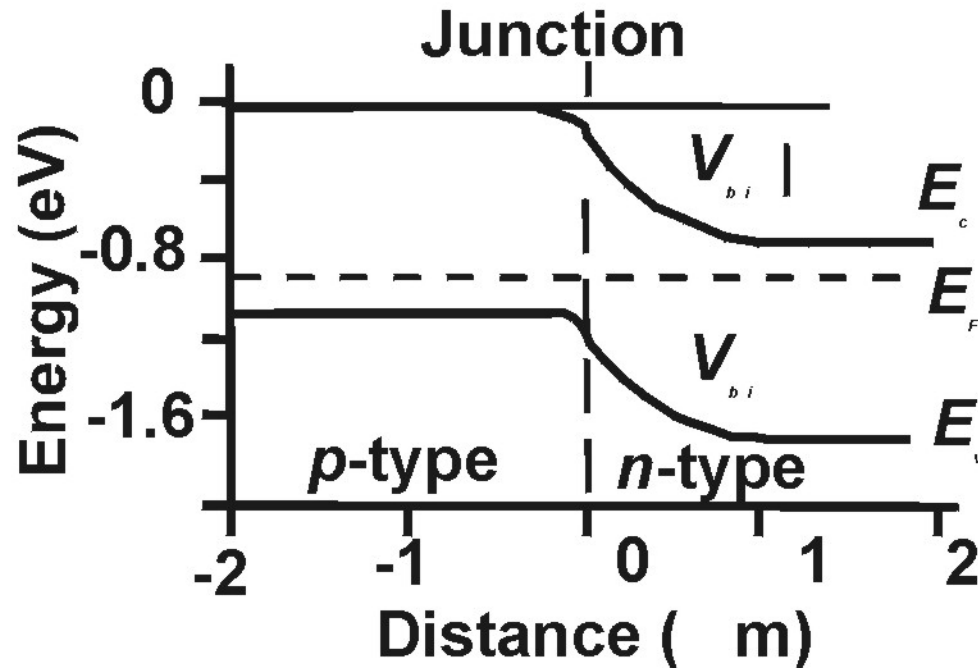
Fermi level,  $E_F$

$$j_p = \mu_p p \frac{\partial E_{Fp}}{\partial x} = \mu_p p \frac{\partial E_F}{\partial x} = 0$$

$$j_n = \mu_n n \frac{\partial E_{Fn}}{\partial x} = \mu_n n \frac{\partial E_F}{\partial x} = 0$$

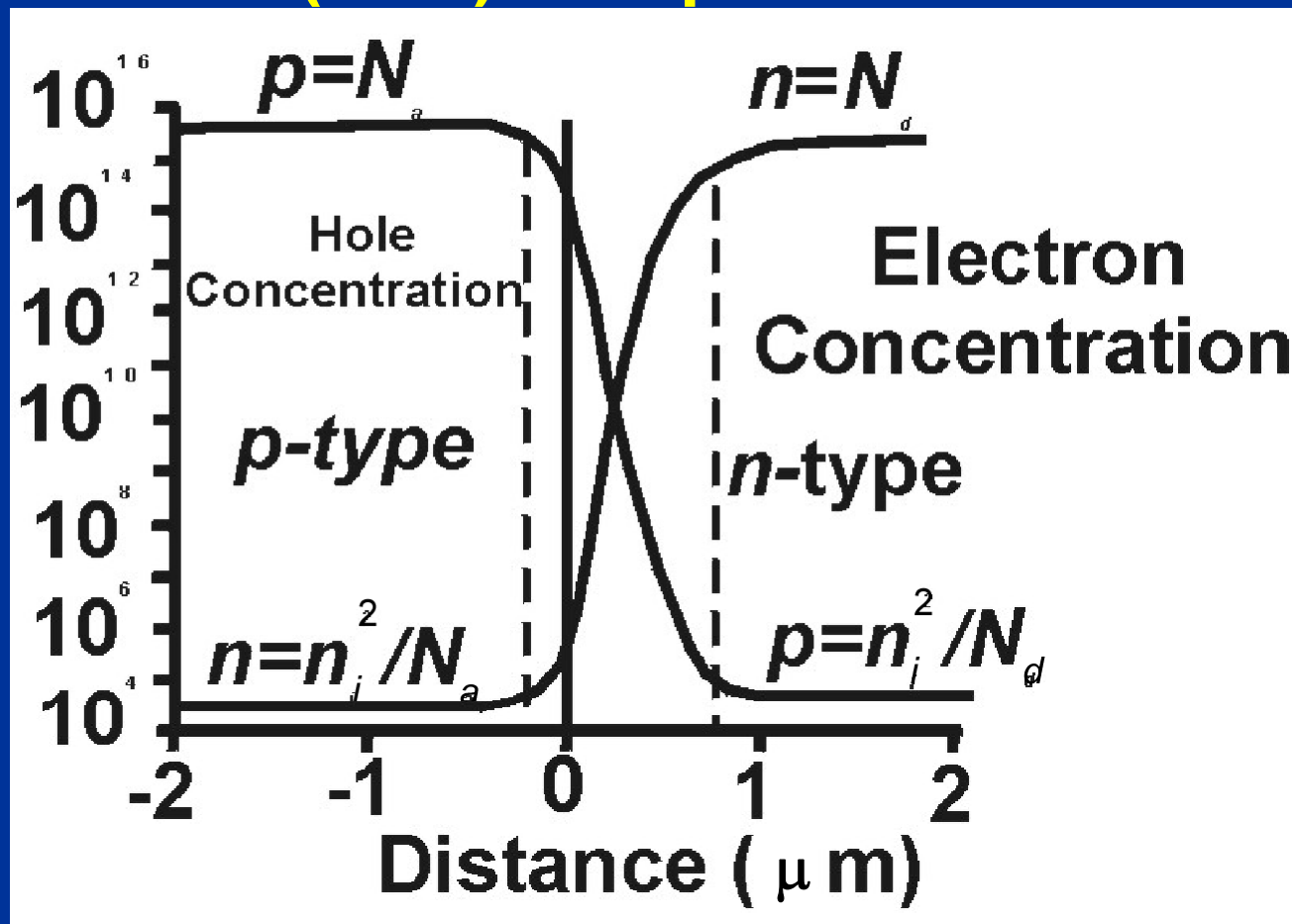
Where  $j_n$ ,  $j_p$ ,  $n$ ,  $p$ ,  $\mu_n$ ,  $\mu_p$ ,  $E_{Fn}$ , and  $E_{Fp}$  are electron and hole current densities, carrier concentrations, mobilities, and quasi-Fermi levels, respectively.

## Energy Band Diagram of Si p-n junction at Equilibrium



Acceptor and donor densities  $5 \times 10^{15} \text{ cm}^{-3}$  and  $10^{15} \text{ cm}^{-3}$ .  $T = 300 \text{ K}$ .

## Electron and Hole Concentrations ( $\text{cm}^{-3}$ ) at Equilibrium



Acceptor and donor densities  $5 \times 10^{15} \text{ cm}^{-3}$  and  $10^{15} \text{ cm}^{-3}$ .  $T = 300 \text{ K}$ .

## Depletion Approximation

Charge density in the *n*-type depletion region

$$\rho_n \approx qN_d$$

Charge density in the *p*-type depletion region

$$\rho_p \approx -qN_a$$

Charge densities of free carriers in the depletion region are much smaller

## Electron and Hole Concentration Profiles in n-type Region

$$n = N_d \exp\left(\frac{\phi(x)}{V_{th}}\right) = N_d \exp\left(-\frac{E_c(x)}{qV_{th}}\right)$$

$$p = N_a \exp\left[-\frac{V_{bi} + \phi(x)}{V_{th}}\right] = N_a \exp\left[\frac{-qV_{bi} + E_c(x)}{qV_{th}}\right]$$

Built-in voltage  $V_{bi} = V_{th} \ln\left(\frac{N_d N_a}{n_i^2}\right)$

## Built-in Voltage

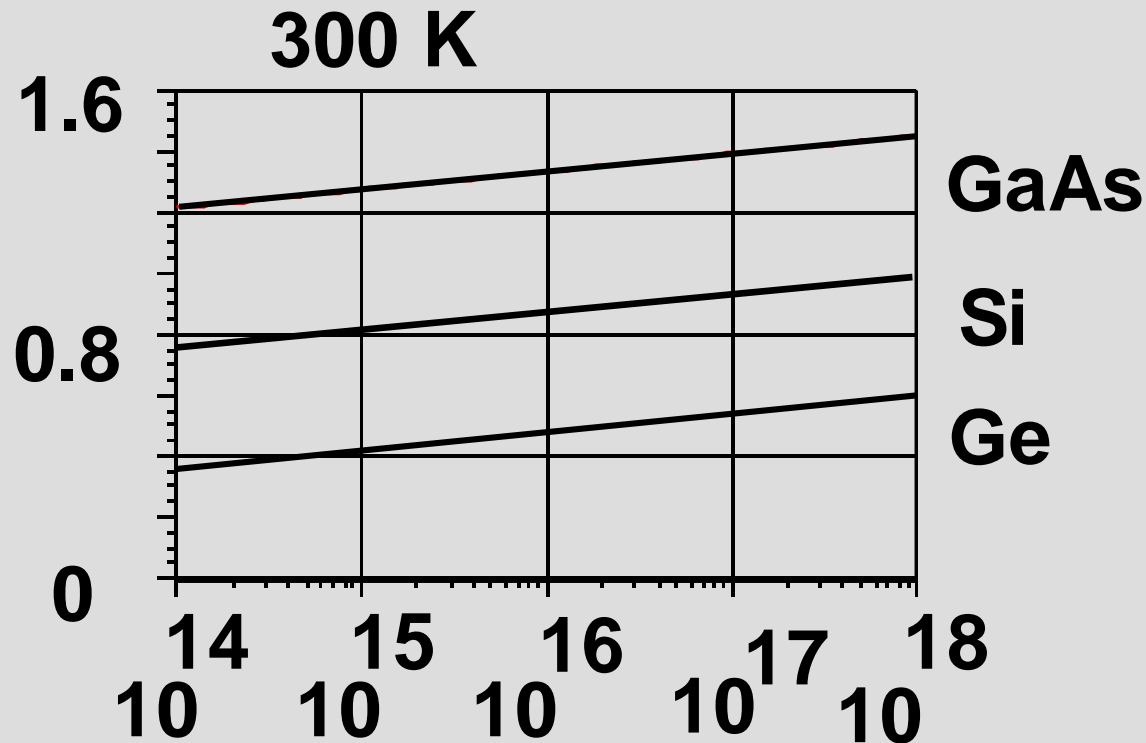
Since,  $n_i = (N_c N_v)^{1/2} \exp\left(-\frac{E_g}{2k_B T}\right)$

$$V_{bi} = E_g + V_{th} \ln\left(\frac{N_d N_a}{N_c N_v}\right)$$

**Built-in voltage is approximately proportional to the energy gap.**

## Built-in Voltage (cont.)

Built-in voltage for n<sup>+</sup>-p junctions  
versus acceptor concentration (cm<sup>-3</sup>)

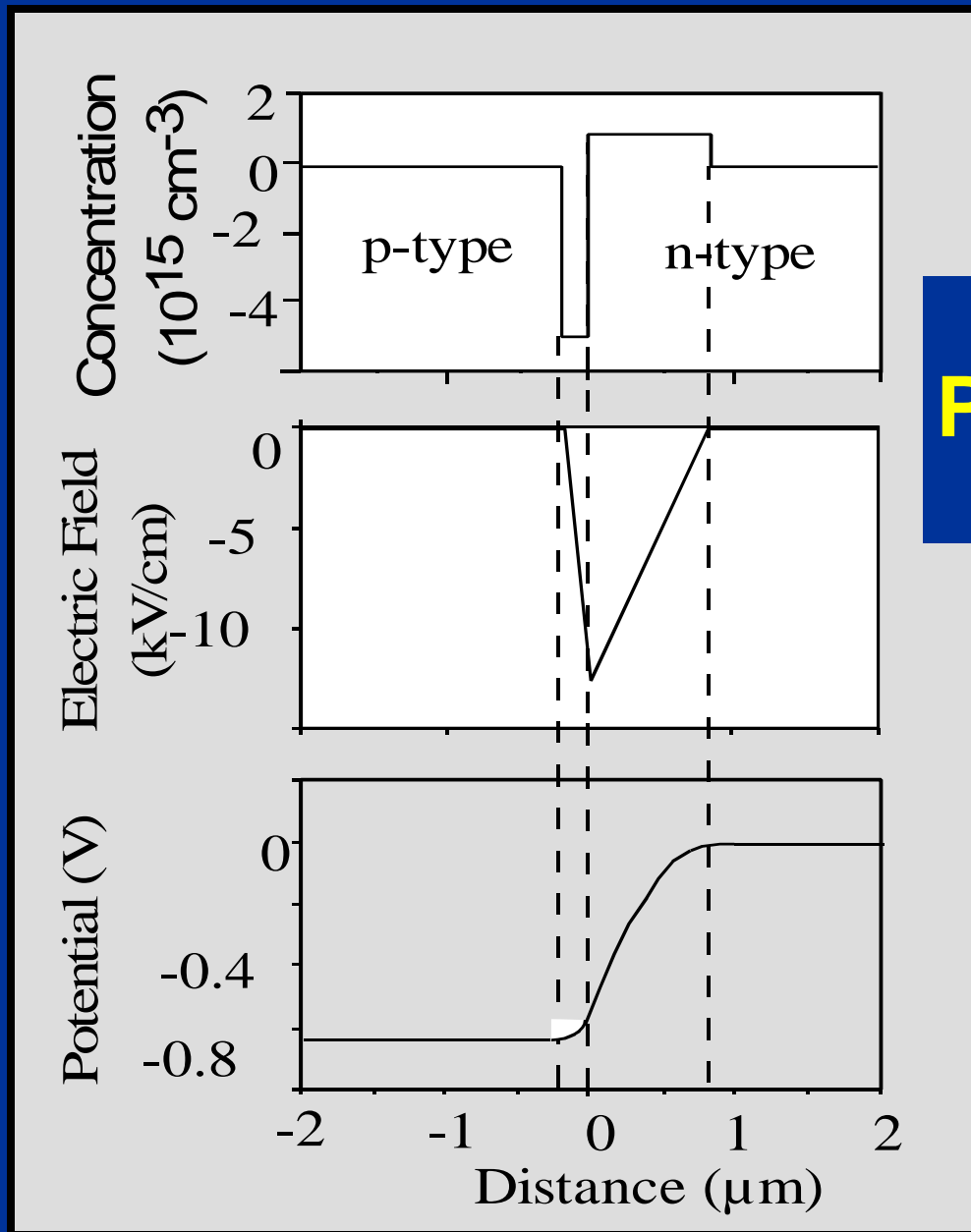


## Charge, Field, Potential Profiles

Using the depletion approximation, we obtain

$$\frac{dF}{dx} = \begin{cases} -\frac{qN_a}{\epsilon_s}, & \text{for } -x_p < x < 0 \\ \frac{qN_d}{\epsilon_s}, & \text{for } 0 < x < x_n \end{cases}$$

$$F = \begin{cases} -F_m \left( 1 + \frac{x}{x_p} \right), & \text{for } -x_p < x < 0 \\ -F_m \left( 1 - \frac{x}{x_n} \right), & \text{for } 0 < x < x_n \end{cases}$$



## Charge, Field, Potential Profiles (cont.)

## Depletion Width

$$\frac{x_n}{x_p} = \frac{N_a}{N_d} \quad \frac{qN_d x_n^2}{2\epsilon_s} + \frac{qN_a x_p^2}{2\epsilon_s} = V_{bi}$$

$$x_n = \sqrt{\frac{2\epsilon_s V_{bi}}{qN_d (1 + N_d / N_a)}}$$

$$x_p = \sqrt{\frac{2\epsilon_s V_{bi}}{qN_a (1 + N_a / N_d)}}$$

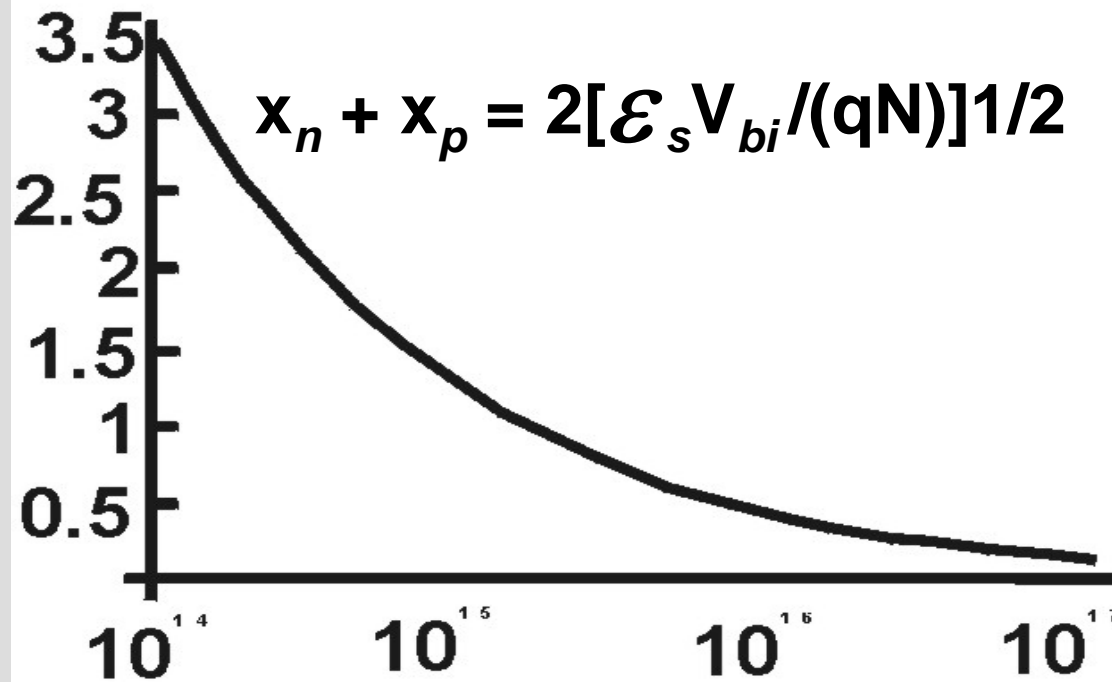
## Example

Calculate and plot the total depletion width,  $x_n + x_p$ , as a function of doping for a Si  $p$ - $n$  junction with the donor concentration in the  $n$ -type region,  $N_d$ , equal to the acceptor concentration in the  $p$ -region,  $N_a$  (such a junction is called a symmetrical junction). Temperature  $T = 300$  K. The dielectric permittivity of Si is  $1.05 \times 10^{-10}$  F/m. The intrinsic carrier concentration is approximately  $10^{10}$  cm<sup>-3</sup>. Vary the doping density from  $10^{14}$  cm<sup>-3</sup> to  $10^{17}$  cm<sup>-3</sup>.

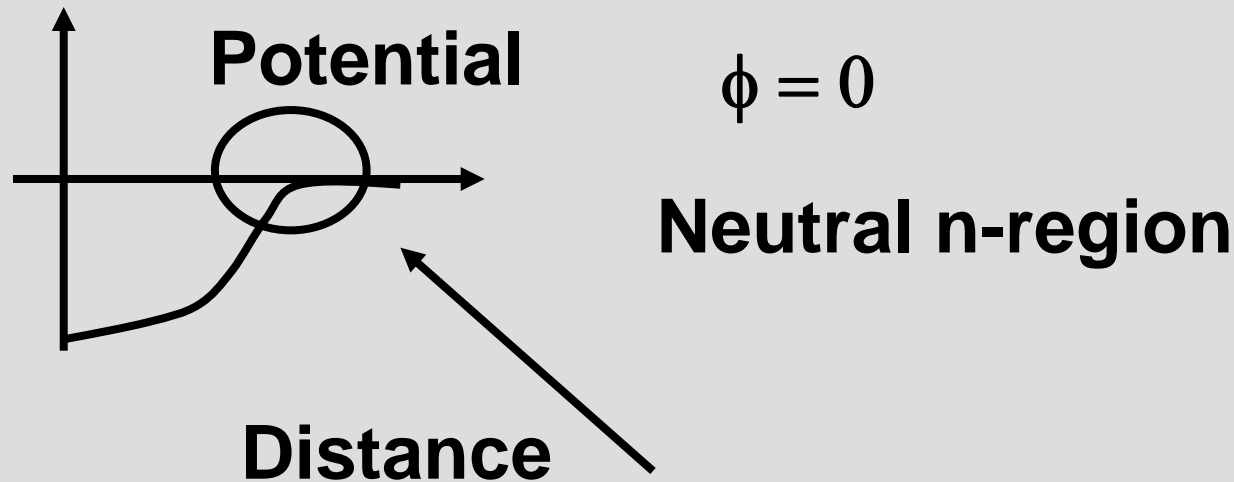
# Solution

For a symmetrical junction,  $V_{bi} = 2V_{th} \ln(N/n_i)$

Depletion  
width  
versus  
doping  
( $\text{cm}^{-3}$ )



## Debye Length - Boundary Region



Let us consider the region where  $\phi < kT/q$

$$\frac{d^2\phi}{dx^2} = -\frac{qN_d}{\epsilon_s} + \frac{qN_d}{\epsilon_s} \exp\left(\frac{\phi}{V_{th}}\right)$$

## Debye Length

Expanding the exponent into a Taylor series

$$\frac{d^2\phi}{dx^2} \approx \frac{qN_d\phi}{\epsilon_s V_{th}} = \frac{\phi}{L_{Dn}^2}$$

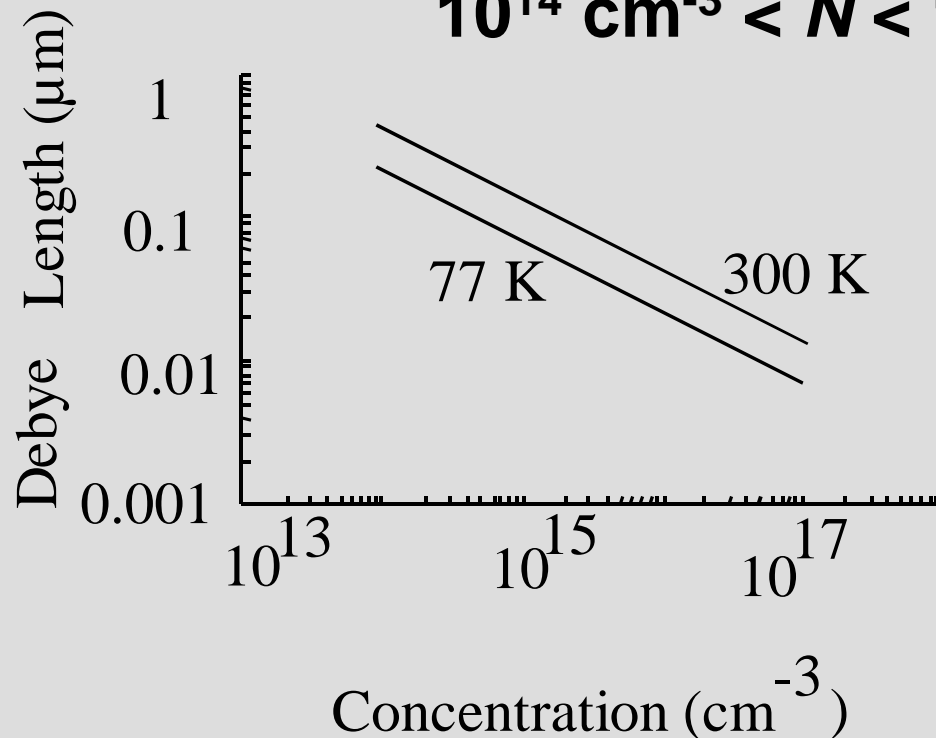
where

$$L_{Dn} = \sqrt{\frac{\epsilon_s V_{th}}{qN_d}}$$

## Example

Calculate and plot the Debye length in GaAs versus doping density at  $T = 77\text{ K}$  and  $300\text{ K}$

$$10^{14}\text{ cm}^{-3} < N < 10^{17}\text{ cm}^{-3}$$



Dielectric permittivity  
 $\epsilon = 1.14 \times 10^{-10}\text{ F/m}$

# Summary of Equations Describing p-n Junctions at Zero Bias

Maximum electric field in a p-n junction	$F_m = \frac{qN_d x_n}{\epsilon_s} = \frac{qN_a x_p}{\epsilon_s}$
Built-in voltage	$V_{bi} = V_{th} \ln(N_d N_a / n_i^2)$
Width of depletion regions	$x_n = \sqrt{\frac{2\epsilon_s V_{bi}}{qN_d(1 + N_d/N_a)}}, \quad x_p = \sqrt{\frac{2\epsilon_s V_{bi}}{qN_a(1 + N_a/N_d)}}$
Debye lengths	$L_{Dn} = \sqrt{\frac{\epsilon_s V_{th}}{qN_d}}, \quad L_{Dp} = \sqrt{\frac{\epsilon_s V_{th}}{qN_a}}$

## p-n JUNCTIONS Under BIAS

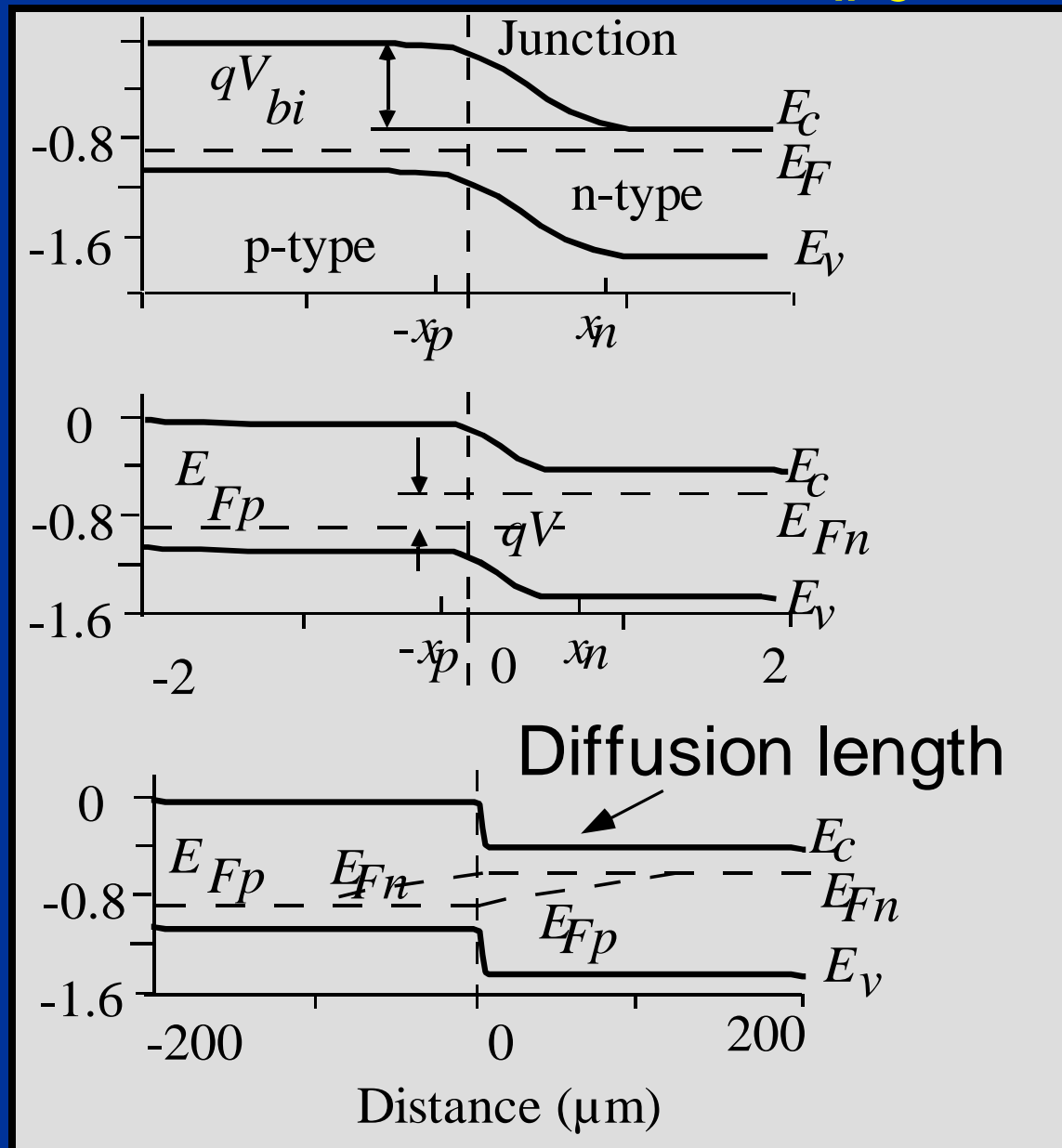
For small voltages, the net current in a p-n junction is still much smaller than the drift and diffusion currents in the depletion region evaluated separately, and, consequently, the electron and hole quasi-Fermi levels defined by

$$n = N_c \exp \frac{E_{Fn} - E_c}{k_B T}$$

$$p = N_v \exp \frac{E_v - E_{Fp}}{k_B T}$$

Remain nearly constant throughout the depletion region.

## Band Diagrams



## Law of the Junction

$$qV = E_{Fn} - E_{Fp}$$
$$pn = n_i^2 \exp\left(\frac{E_{Fn} - E_{Fp}}{k_B T}\right) = n_i^2 \exp\left(\frac{V}{V_{th}}\right)$$

### Example

The intrinsic carrier concentration in GaAs at 300 K is approximately  $10^6 \text{ cm}^{-3}$ . Estimate the forward bias required to create the electron-hole density of  $10^{17} \text{ cm}^{-3}$  in the depletion region at the point where  $p = n$ .

## Solution

From the law of the junction

$$V = V_{th} \ln \frac{pn}{v_i^2} = 0.02584 \times \ln \frac{10^{34}}{10^{12}} = 1.308 V$$

## Diffusion Length of Minority Carriers

Diffusion Equation  $D_p \frac{\partial^2 p_n}{\partial x^2} - \frac{p_n - p_{no}}{\tau_{pl}} = 0$

Solution of the Diffusion Equation

$$p_n(x) - p_{no} = A \exp\left(\frac{x - x_n}{L_p}\right) + B \exp\left(-\frac{x - x_n}{L_p}\right)$$

Hole Diffusion Length  $L_p = \sqrt{D_p \tau_{pl}}$

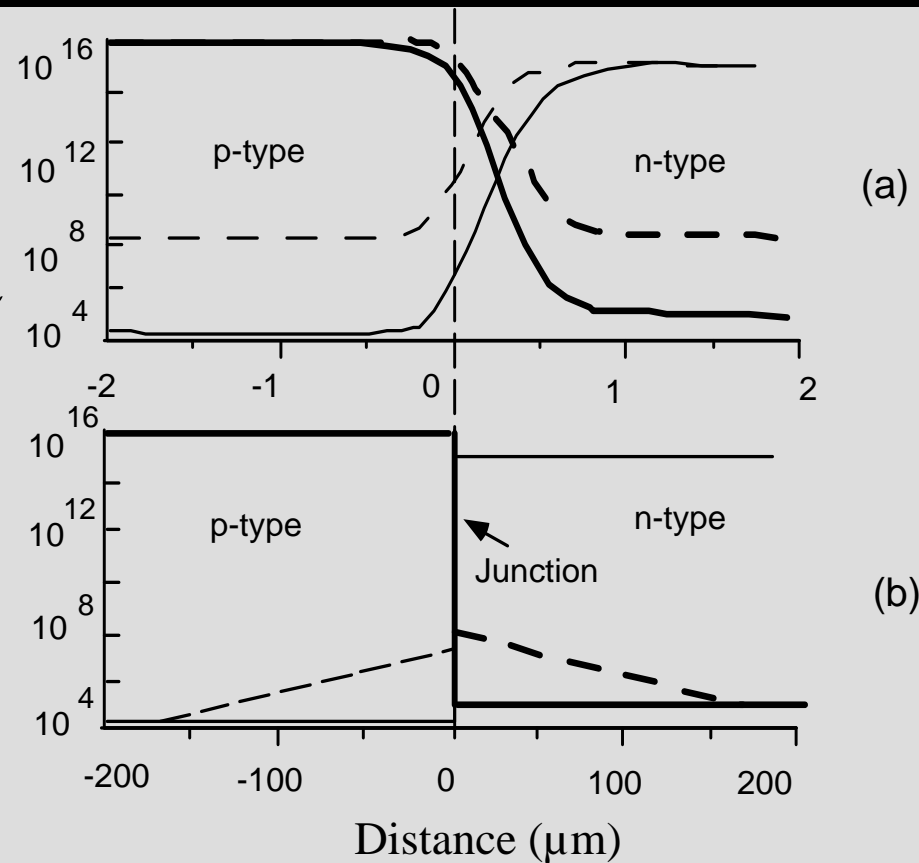
# Concentration Profiles

**Boundary conditions:**

$$p_n(x_n) = p_{no} \exp\left(\frac{V}{V_{th}}\right)$$

$$p_n(x \rightarrow \infty) = p_{no}$$

Concentration (cm<sup>-3</sup>)



**Solution:**

$$p_n(x) - p_{no} = p_{no} \left[ \exp\left(\frac{V}{V_{th}}\right) - 1 \right] \exp\left(-\frac{x - x_n}{L_p}\right)$$

## Example

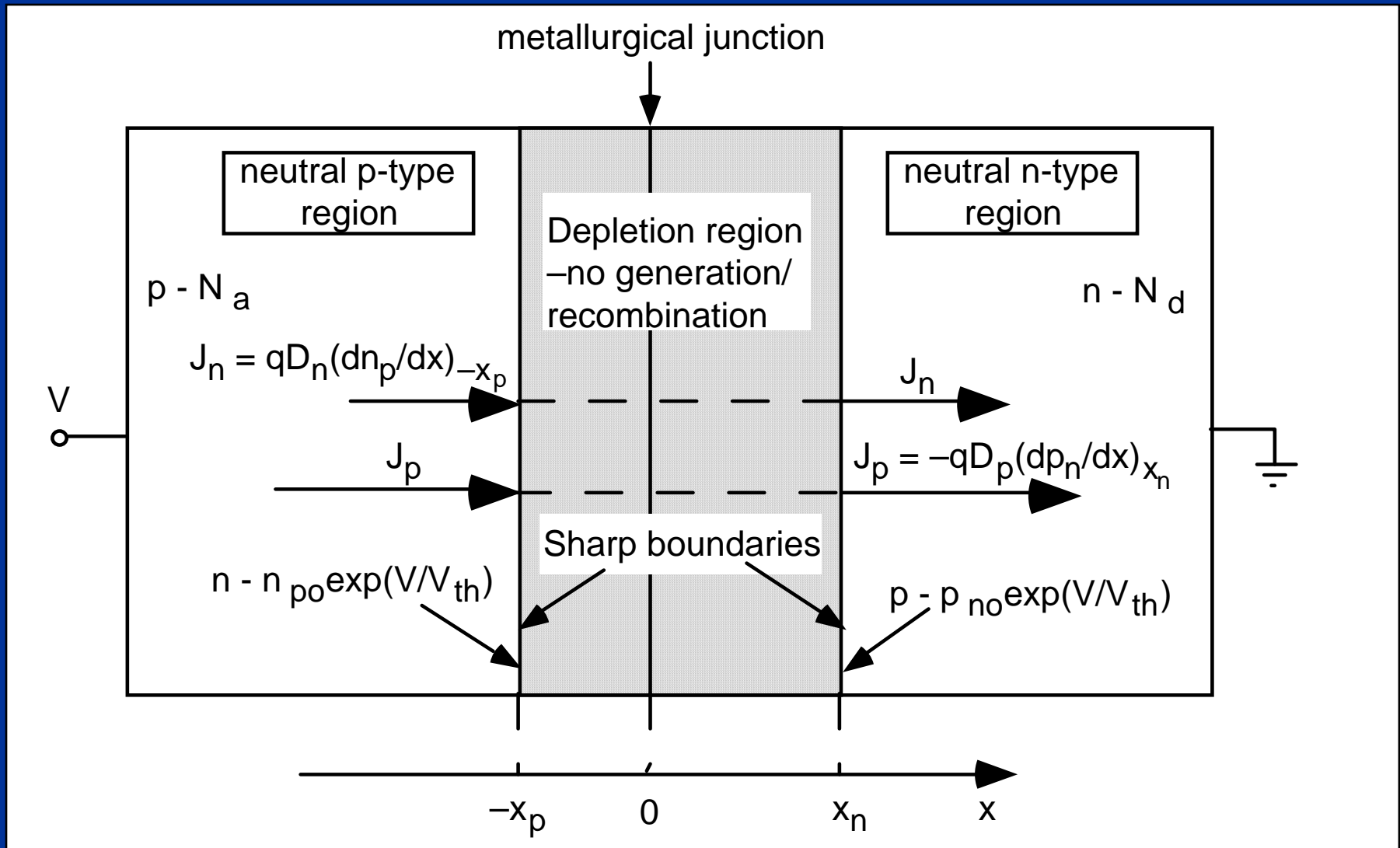
The slope of the semilog dependence of the hole concentration on a distance in the  $n$ -type region of a  $p$ - $n$  diode at room temperature ( $T = 300$  K) is 1 decade/ $120\mu\text{m}$ . The hole mobility is  $250\text{ cm}^2/\text{V}\cdot\text{s}$ . What is the hole lifetime?

## Solution

The slope in the semilog scale is equal to  $1/[L_p \ln(10)]$ . Hence  $L_p = 120/\ln(10) = 52.1$  m. Using the Einstein relation, we find  $D_p = V_{th} = 250 \times 0.02584 = 6.46$  (cm<sup>2</sup>/s). Hence, the hole lifetime,

$$\tau_{pl} = \frac{L_p^2}{D_p} = \frac{(52.1 \times 10^{-4})^2}{6.46} = 4.20 \times 10^{-6} \text{ (s)}$$

# I-V Characteristics of p-n Junctions



## Assumptions

- All applied voltage drops across depletion region
- Low-level injection only
- Sharp depletion boundaries
- No generation-recombination processes inside depletion region

## Procedure

- Find minority carrier concentration at both boundaries (Use the Law of the Junction)
- Find distribution of minority carriers inside neutral regions (Diffusion equation)
- Find minority carrier diffusion current at boundaries (Gradients of concentrations)
- Total current density is the sum of the minority carrier diffusion currents at the two boundaries ( $J = J_n + J_p$ )

# For the Ideal p-n Junction

The ideal diode equation:

$$J = J_s \left[ \exp\left(\frac{V}{V_{th}}\right) - 1 \right]$$

Saturation current:

$$J_s = \frac{qD_p p_{n0}}{L_p} + \frac{qD_n n_{p0}}{L_n}$$

Diffusion lengths:

$$L_n = \sqrt{D_n \tau_n} \quad L_p = \sqrt{D_p \tau_p}$$

## Ideal Diode Equation

$$j_p \approx j_{pD} = -qD_p \frac{\partial p_n}{\partial x} = \frac{qD_p p_{no}}{L_p} \left[ \exp\left(\frac{V}{V_{th}}\right) - 1 \right] \exp\left(-\frac{x - x_n}{L_p}\right)$$

$$j \approx j_{Dp} \Big|_{x=x_n} + j_{Dn} \Big|_{x=-x_p}$$

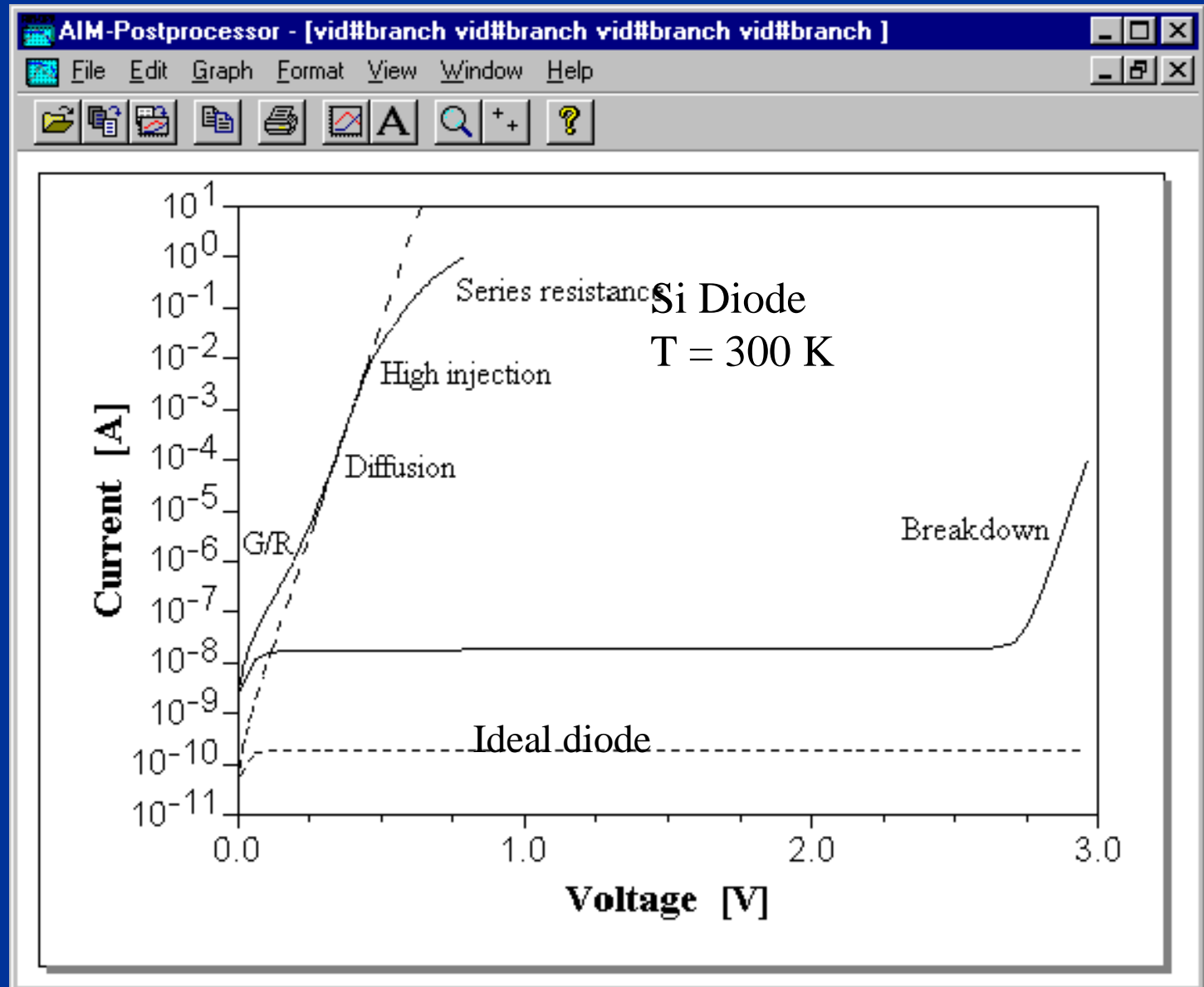
$$\left| I = I_s \left[ \exp\left(\frac{V}{V_{th}}\right) - 1 \right] \right| \quad I = j S$$

The diode saturation current is

$$I_s = S \left( \frac{qD_p p_{no}}{L_p} + \frac{qD_n n_{po}}{L_n} \right)$$

## Non-Ideal Diode Effects

From T. Fjeldly,  
T. Ytterdal,  
and M. S. Shur,  
Introduction to device  
modeling and circuit  
simulation, Wiley (1997)



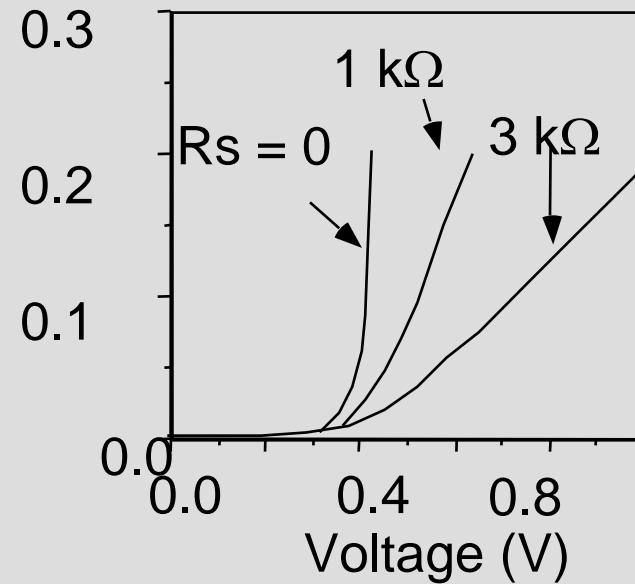
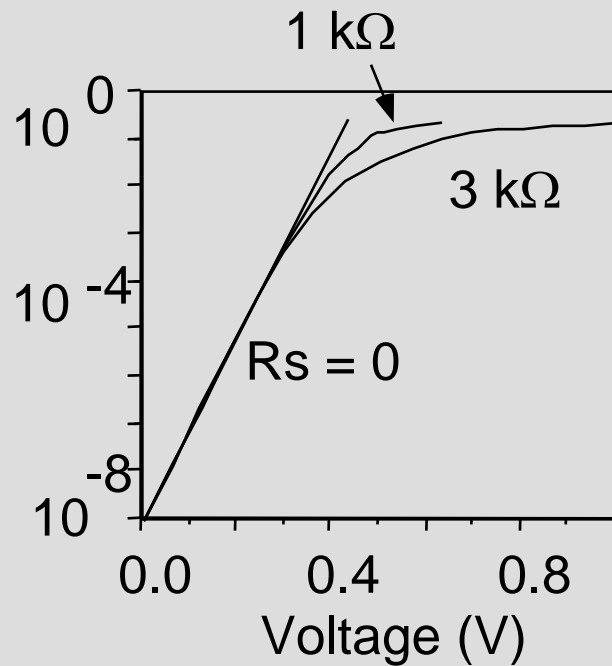
## A. Series Resistance

$$I = I_s \left[ \exp\left(\frac{V - IR_s}{V_{th}}\right) - 1 \right]$$

## B. High Injection

- At high forward bias ( $V \Rightarrow V_{bi}$ ), the injected minority carrier concentration becomes comparable to the majority carrier concentration
- The neutrality condition demands a similar increase in both electron and hole concentrations
- This reduces the growth in the minority carrier concentration and in the current

## Effect of Series Resistance



**Current-voltage characteristics of p-n diode**  
 $I_{\text{seff}} = 10^{-8}$  mA. The characteristics are shown in semilog and linear scales. (Assumed  $\eta = 1$ )

## B. High Injection

Law of the junction: 
$$\begin{cases} pn = p(N_d + p) = n_i^2 \exp(V/V_{th}) \\ \Rightarrow p \rightarrow n_i \exp(V/2V_{th}) \text{ when } p \geq N_d \end{cases}$$

Diode current: 
$$I \propto \exp(V/2V_{th})$$

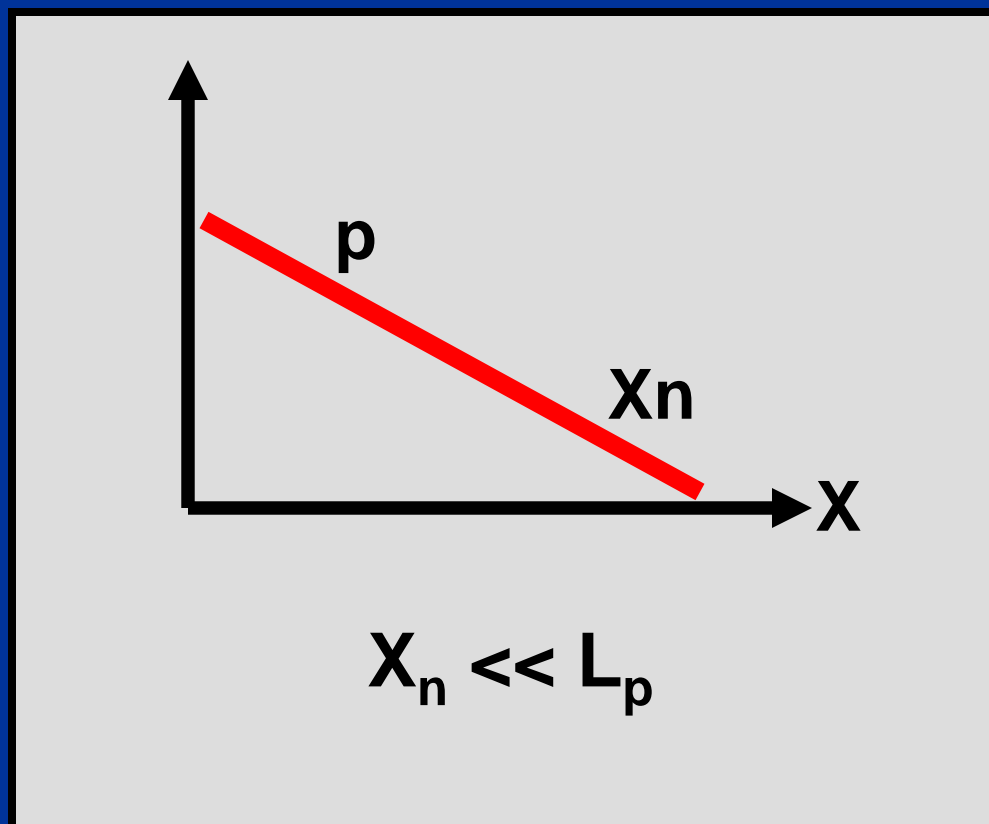
## C. G/R Processes in the Depletion Region

*G/R processes seek to restore thermal equilibrium:  $pn \Rightarrow n_i^2$*

Reverse biased  $pn < n_i^2$

$$J_{gen} = q \int_{-x_p}^{x_n} |G - R| dx = \frac{qn_i x_d}{\tau_{gen}}$$

## Short p<sup>+</sup>-n Diode



## Current-Voltage Characteristic

When the length of the  $n$ -section,  $X_n$ , is much smaller than the hole diffusion length  $L_p$ , the solution of the diffusion equation is given by

$$p_n(x) \approx p_{no} + p_{no} \left[ \exp\left(\frac{V}{V_{th}}\right) - 1 \right] \left( \frac{X_n - x}{X_n - x_n} \right)$$

and the hole distribution in the  $n$ -type region is a linear function of  $x$ . In this case, the current density is given by

$$j \approx -qD_p \left. \frac{\partial p_n}{\partial x} \right|_{x=x_n} = \frac{qD_p p_{no}}{X_n - x_n} \left[ \exp\left(\frac{V}{V_{th}}\right) - 1 \right]$$

This current density is  $L_p/(X_n - x_n)$  larger than for a long  $p^+ - n$  diode

## Generation Current

Under thermal equilibrium, the thermal generation of electron-hole pairs is balanced by their recombination ( $G = R$ ). Under a reverse bias, there are very few electron-hole pairs in the depletion region, and the recombination rate,  $R$ , is nearly zero. However, the thermal generation processes continuously supply the electron hole pairs into the depletion region. The thermal rate can be estimated as

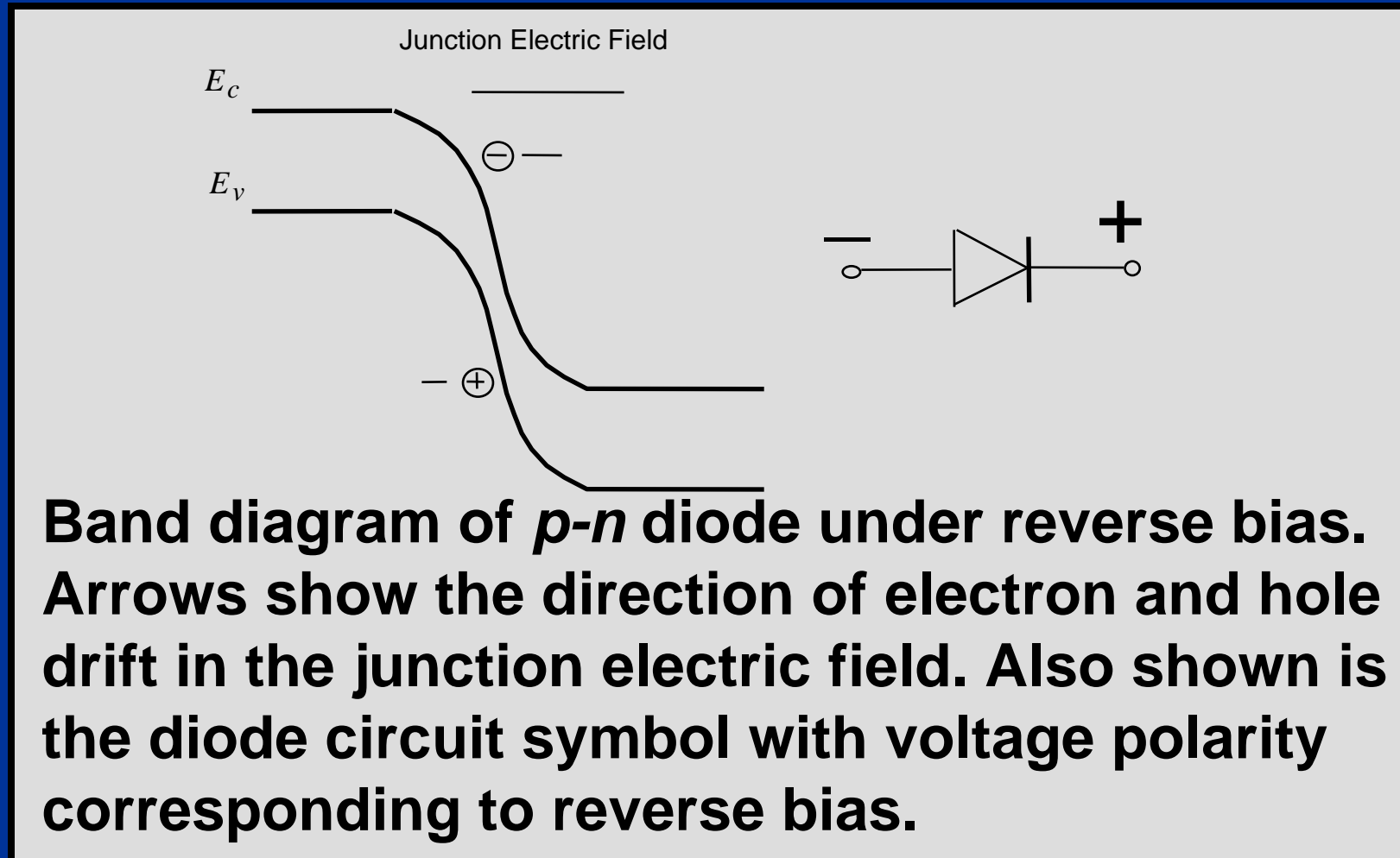
$$G_{thermal} = \frac{n_i}{\tau_{gen}}$$
 where  $\tau_{gen}$  is the effective generation time of electron hole pairs in depletion region.

## Generation Current (cont.)

The total electron charge (equal to the hole charge) supplied into the depletion region per unit area per second is equal to

$qG_{thermal}x_d = qn_i x / \tau_{gen}$  where  $x_d = x_n + x_p$  is the width of the depletion region. Since the generation times typically vary between a microsecond and a nanosecond, orders of magnitude higher than  $t_{tr}$ , the generated carriers are swept away almost instantaneously leading to the generation current density

## Electrons and Holes Swept by Reverse Bias



## Diffusion and Generation Currents

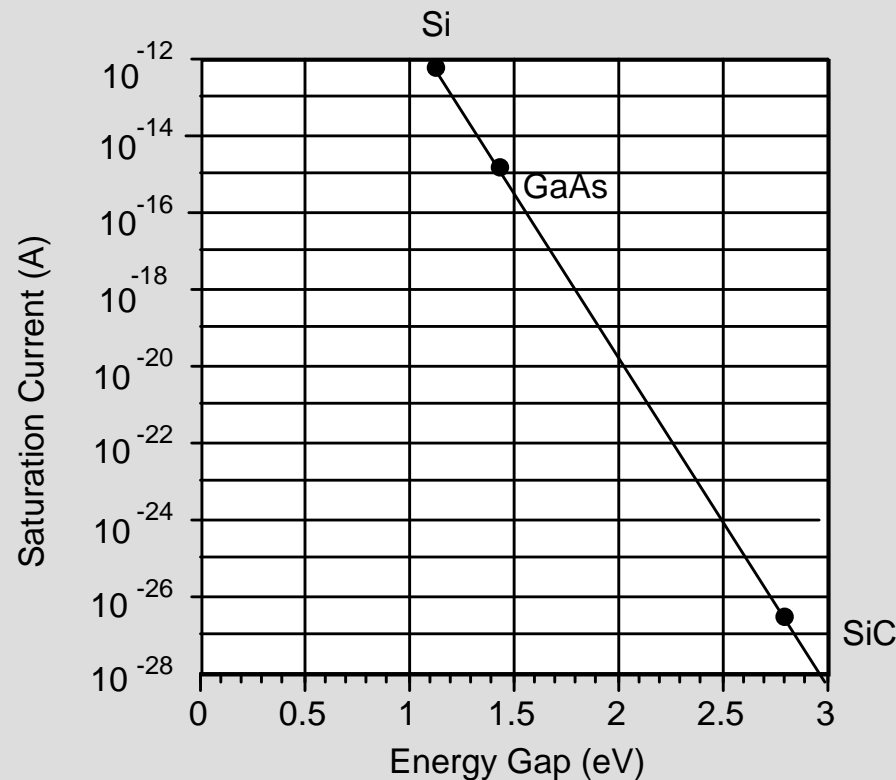
The total reverse current density,  $j_R$ , is given by:

$j_R = j_s + j_{gen}$  where the diffusion component

$$j_s = \left( \frac{qD_p}{N_d L_p} + \frac{qD_n}{N_a L_n} \right) n_i^2 \quad j_{gen} = \frac{qn_i X_d}{\tau_{gen}}$$

Since  $j_s$  is proportional to  $n_i^2$  and  $j_{gen}$  is proportional to  $n_i$  is sufficiently small. In practice, this is often the case for Si, GaAs, and wider gap semiconductors at room temperature and lower temperatures.

## Generation Current in Different Semiconductors



**Generation current (per  $1 \mu\text{m}^3$  of the depletion region volume) versus energy gap (at room temperature).**

## Recombination Current

Under forward bias conditions ( $V > 0$ ) excess electrons and holes are injected into the depletion region where some of them recombine. The recombination current density is equal to the total electron charge per unit area recombining in the depletion region in one second:

$$j_{rec} = q \int_{-x_p}^{x_n} U_R dx$$

## Recombination Current (cont.)

(This electron charge is equal to the total hole charge per unit area recombining in the depletion region.) Here  $U_R$  is the net recombination rate. For a simple model accounting only for one impurity (trap) energy level near the middle of the energy gap, this integral can be evaluated, leading to the following expression:

$$j_{rec} = j_{recs} \exp\left(\frac{V}{2V_{th}}\right) \text{ where } j_{recs} \approx \frac{\pi}{2} \frac{qn_i V_{th}}{\tau_{rec} F_{max}}$$

In practical devices ,  $j_{rec} = j_{recs} \exp\left(\frac{V}{m_r V_{th}}\right)$

Here  $m_r$  may differ from two.

## Empirical Diode Equation

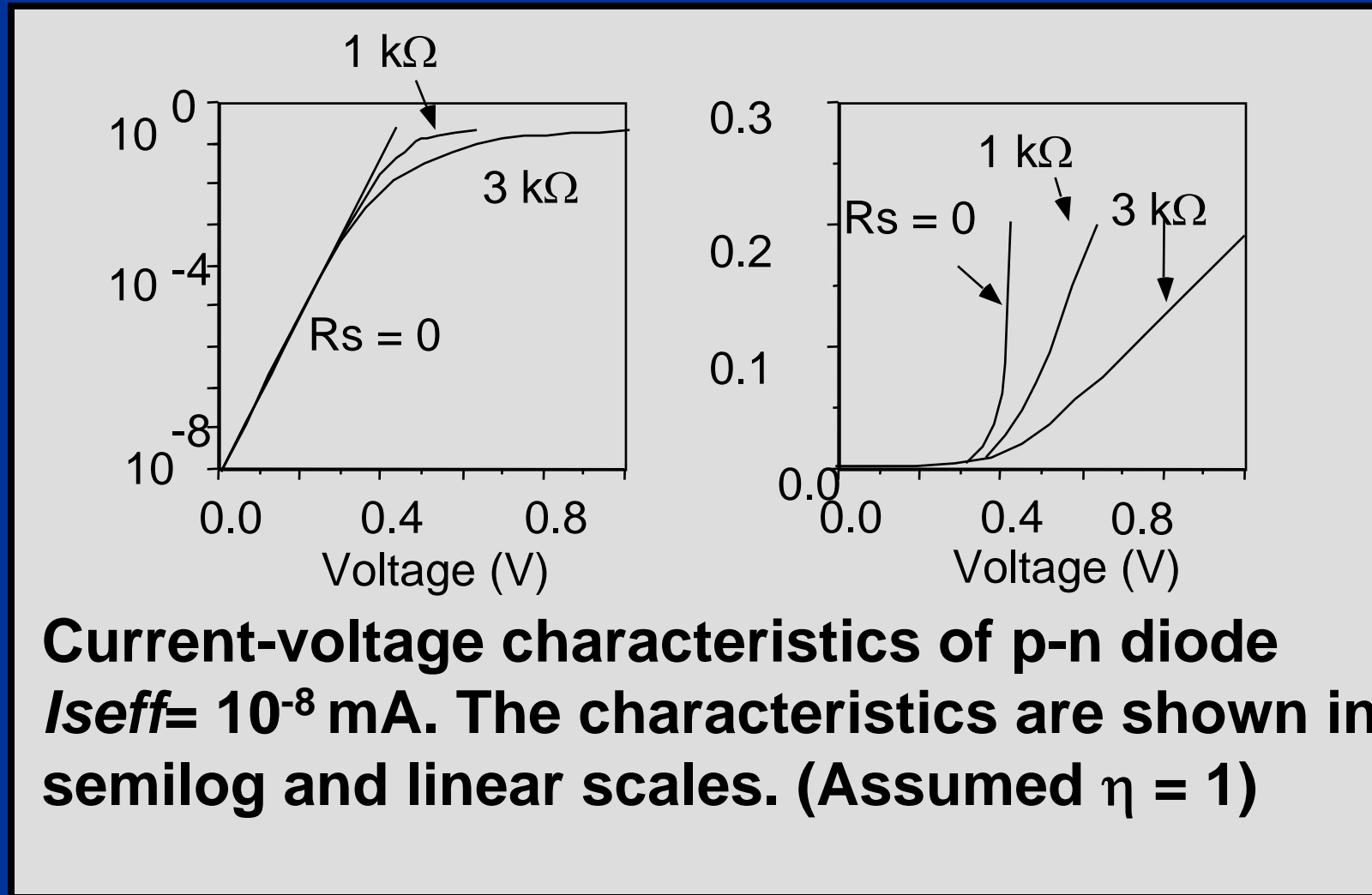
In a real semiconductor diode, the parasitic series resistance,  $R_s$ , of the device contacts and the semiconductor neutral regions may play an important role, and this equation has to be modified to include the voltage drop across the series resistance:

$$I = I_s \left[ \exp\left(\frac{V - IR_s}{\eta V_{th}}\right) - 1 \right]$$

or

$$V = IR_s + \eta V_{th} \ln\left(\frac{I}{I_s}\right)$$

## Effect of Series Resistance (repeat again)



**Current-voltage characteristics of p-n diode**  
 $I_{seff} = 10^{-8}$  mA. The characteristics are shown in semilog and linear scales. (Assumed  $\eta = 1$ )

## Diode Model in SPICE

$$I_s(T) = I_s(T_0) \left( \frac{T}{T_0} \right)^{\frac{k}{\eta}} \exp\left( \frac{E_g}{k_B T_0} \right) \exp\left( - \frac{E_g}{k_B T} \right)$$

Here  $E_g$  is called activation energy (with a default value set to be equal to the energy gap),  $T$  is the device temperature  $T_0$ , is a nominal device temperature (at which the device parameters are specified in the circuit simulator; usually  $T_0 = 300$  K),  $k$  is an empirical temperature exponent,  $k=3$  and  $h=1$  for  $p-n$  diodes when the diffusion saturation current is dominant.

$$I_{leakage} = G_{min} V \quad \text{Parasitic leakage current}$$

## Example

The default value of  $G_{min}$  in SPICE is  $10^{-12} 1/\Omega$ . The doping level of the  $n$ -region of a Si  $p^+n$  diode  $N_d = 10^{15} \text{ cm}^{-3}$ , the generation time  $\tau_{gen} = 10^{-8} \text{ s}$ , the intrinsic carrier density  $n_i = 10^{10} \text{ cm}^{-3}$ , the hole diffusion length  $L_p = 100 \mu\text{m}$ , the hole diffusion coefficient  $D_p = 10 \text{ cm}^2/\text{V}$ , the length of the neutral  $n$ -type region  $X_n - x_n$  is  $10 \mu\text{m}$ , the dielectric permittivity  $\epsilon_s = 1.05 \times 10^{-10} \text{ F/m}$ , the built-in voltage  $V_{bi} = 0.6 \text{ V}$ , the diode cross section  $S = 10^{-2} \text{ cm}^2$ . Compare the saturation diffusion current, generation current and parasitic leakage current at  $10 \text{ V}$  reverse bias (using  $G_{min} = 10^{-12} 1/\Omega$ )

## Solution Part 1

$$I_{leakage} = 10^{-12} \left( \frac{1}{\Omega} \right) \times 10(V) = 10^{-11} A$$

$$\text{Since } p_{no} = \frac{n_i^2}{N_d} = \frac{10^{20}}{10^{15}} = 10^5 (cm^{-3}) \gg n^{po},$$

and  $X_n - x_n \ll L_p$ , using eq. (4-3-21) and converting to the SI units, we find

$$I_{Ds} \approx \frac{qD_p p_{no} S}{X_n - x_n} = \frac{1.602 \times 10^{-19} \times 10 \times 10^{-4} \times 10^{11} \times 10^{-6}}{10 \times 10^{-6}}$$

$$= 1.602 \times 10^{-12} (A)$$

## Solution (Part 2)

For a  $p^+-n$  diode from eqs. (4-3-24) and (4-3-25), we find

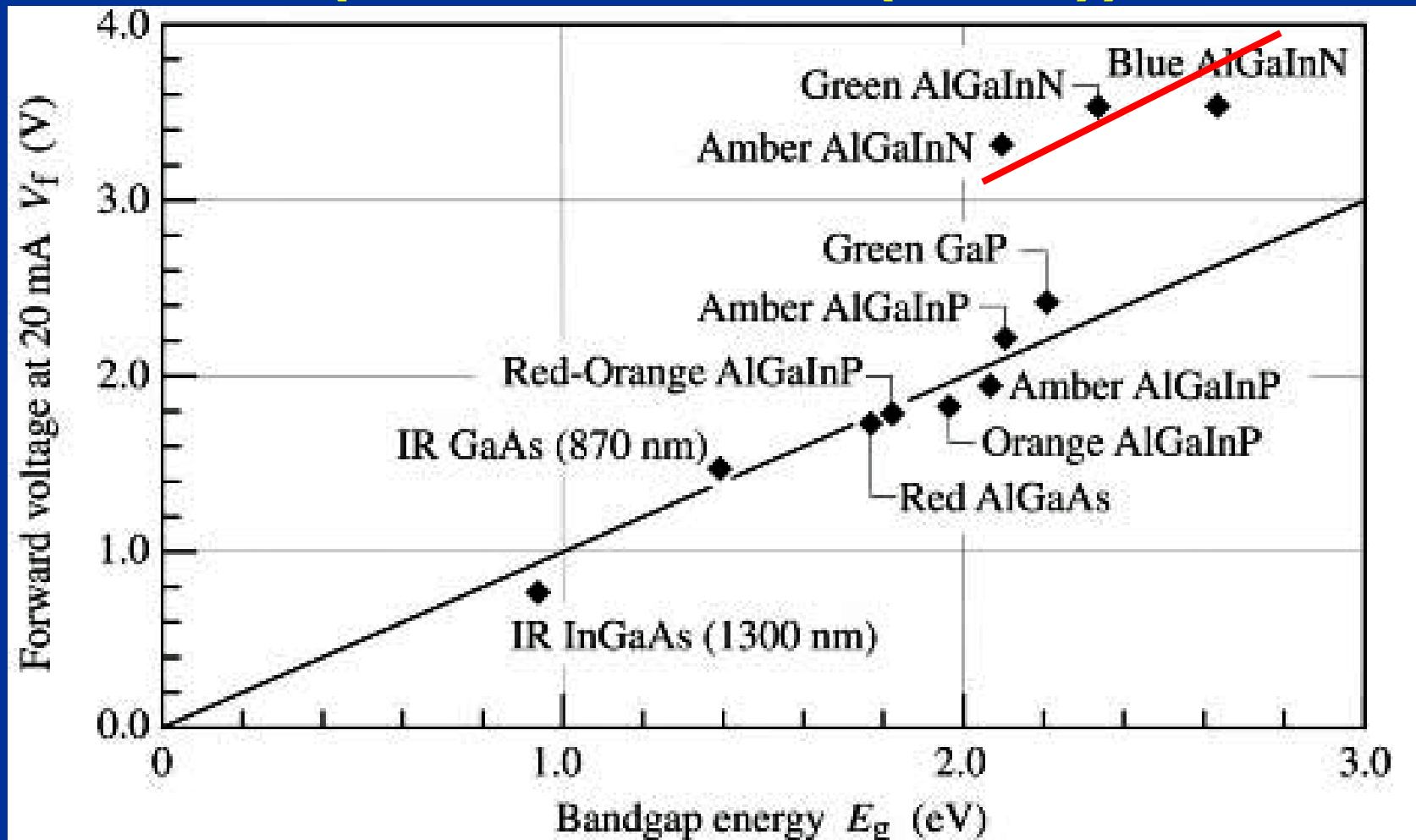
$$x_n = \sqrt{\frac{2\varepsilon_s(V_{bi} - V)}{qN_d}} = \sqrt{\frac{2 \times 1.05 \times 10^{-10} \times [0.6 - (-10)]}{1.602 \times 10^{-19} \times 10^{21}}}$$

$$I_{gen} = \frac{qn_i x_d S}{\tau_{gen}} = 1$$

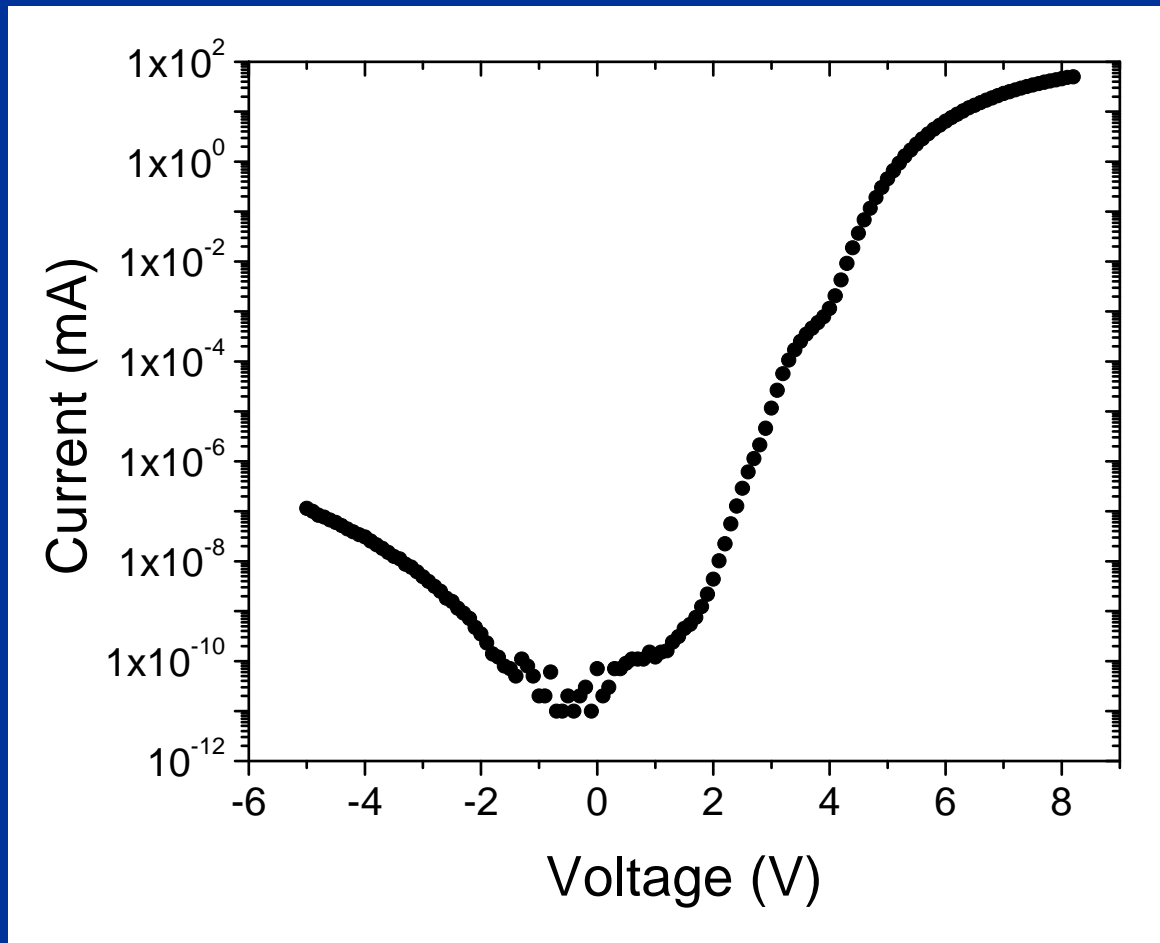
$$\frac{602 \times 10^{-19} \times 10^{16} \times 3.73 \times 10^{-6} \times 10^{-6}}{10^{-8}} = 5.97 \times 10^{-7} \text{ (A)}$$

In this example, the leakage current determined by the default SPICE parameter is smaller. The generation current here is dominant

# Forward LED Voltages (after Krames (2000))

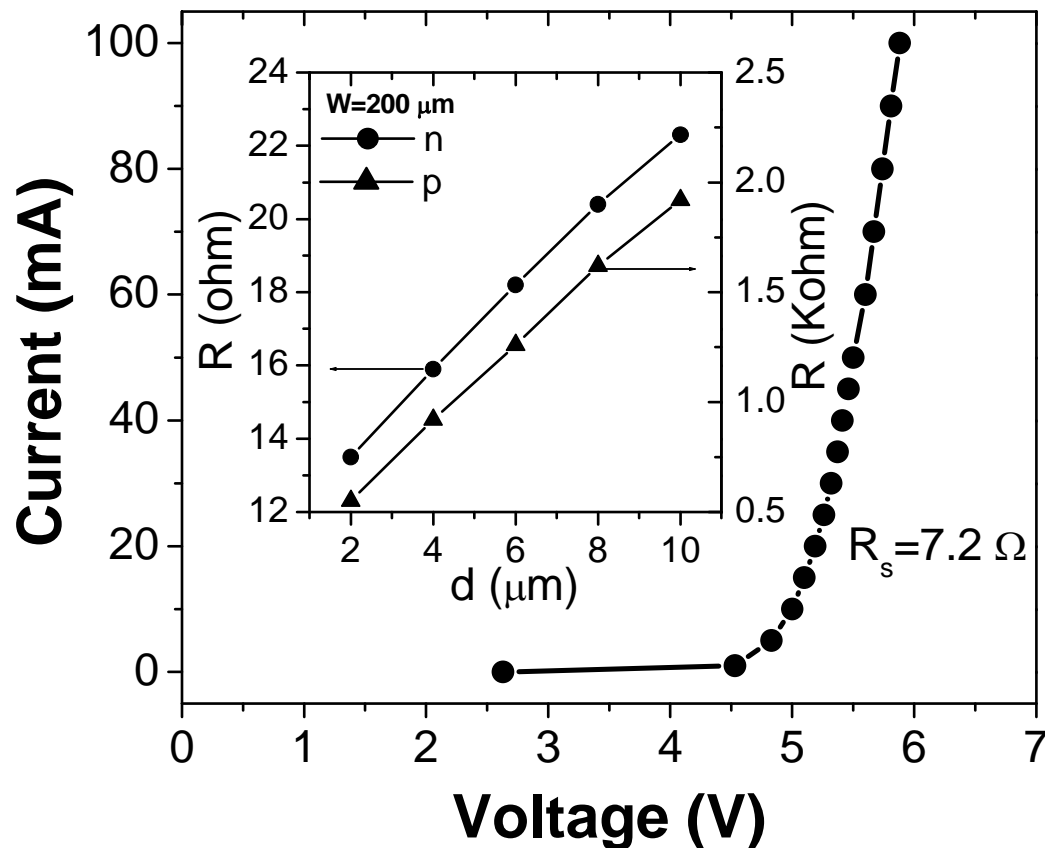


## IV characteristic for a $100 \times 100 \mu\text{m}^2$ 265 nm LED



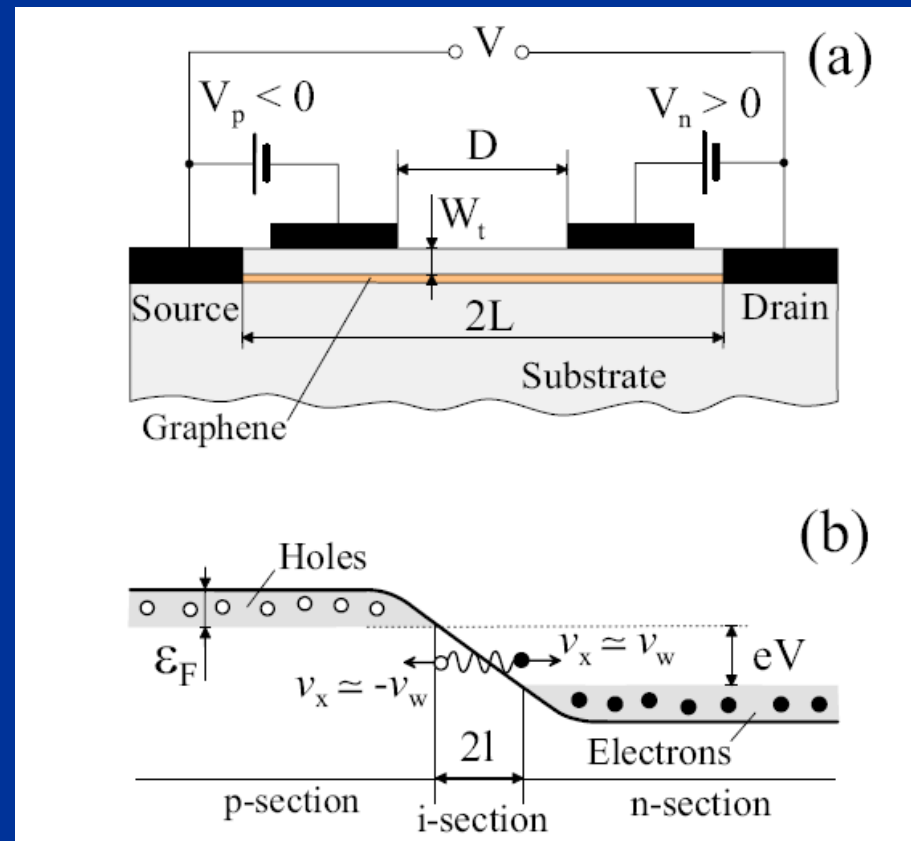
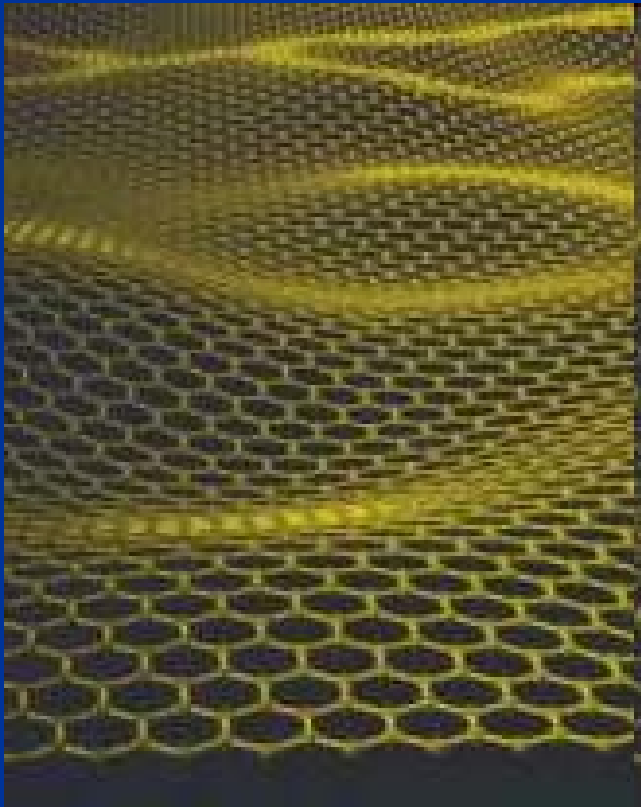
From Y. Bilenko, A. Lunev, X. Hu, J. Deng, T. M Katona, J. Zhang, R. Gaska, M. S Shur, W. Sun, V. Adivarahan, M. Shatalov, and A. Khan, JJAP, Express Letter, Vol. 44, No. 3, pp. L98-L100 (2005)

# IV curve for the two $0.0001 \text{ cm}^2$ 280 nm parallel packaged UV LEDs. Inset showing the p- and n-TLM results



From J. P. Zhang, X. Hu, Y. Bilenko, J. Deng, A. Lunev, M. S. Shur, R. Gaska, M. Shatalov, J. W. Yang, M. A. Khan, Appl. Phys. Lett., Vol. 85, pp. 5532-5534, No 23, 2004

# Graphene p-n junction



From V. Ryzhii, M. Ryzhii, V. Mitin and M. S. Shur Graphene Tunneling Transit-Time Terahertz Oscillator Based on Electrically Induced p-i-n Junction

From <http://images.iop.org/objects/physicsweb/news/thumb/12/3/23/Graphene.jpg>