

## Chapter 5-1. PN-junction electrostatics

In this chapter you will learn about pn junction electrostatics:  
Charge density, electric field and electrostatic potential existing  
inside the diode under equilibrium and steady state conditions.

**You will also learn about:**

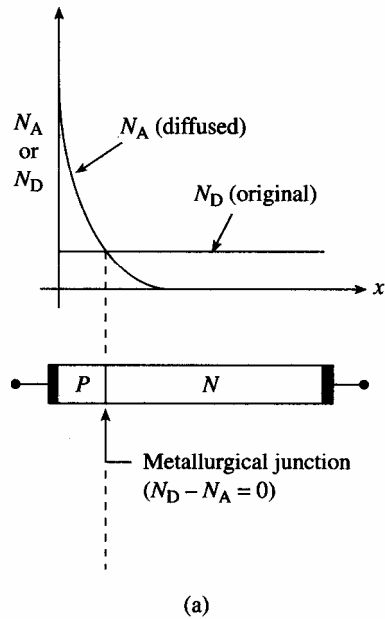
Poisson's Equation

Built-In Potential

Depletion Approximation

Step-Junction Solution

# PN-junction fabrication



PN-junctions are created by several processes including:

1. Diffusion
2. Ion-implantation
3. Epitaxial deposition

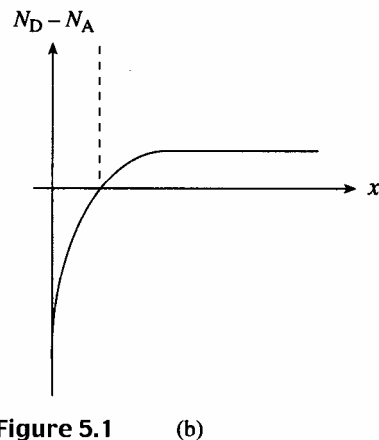
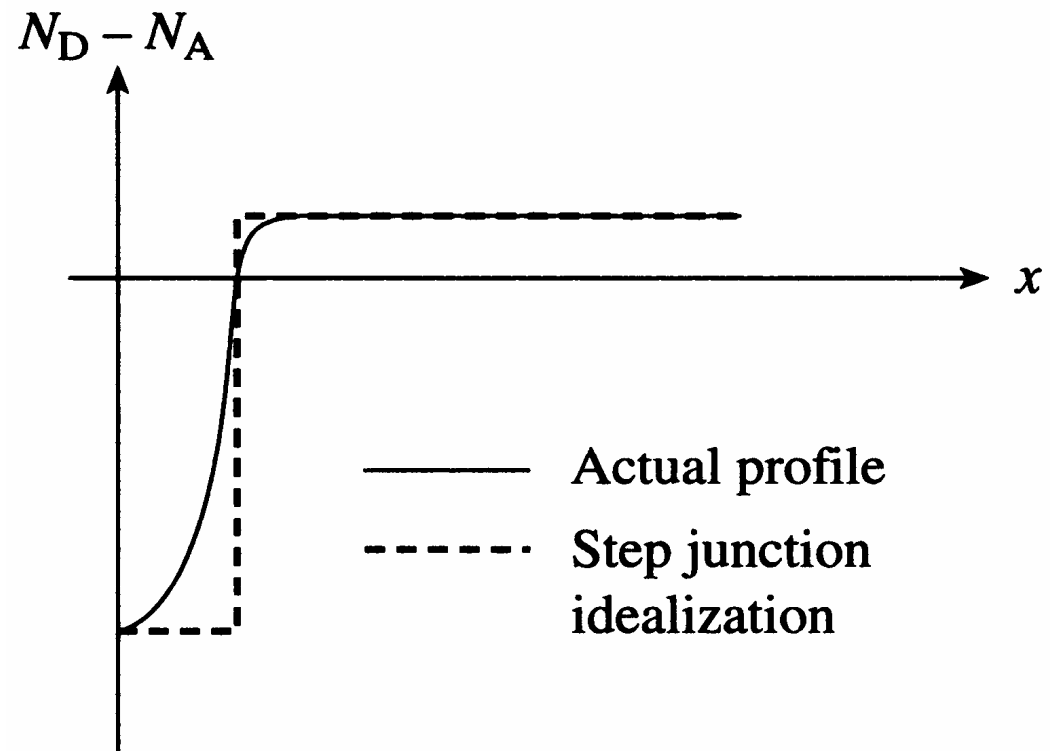


Figure 5.1 (b)

Each process results in different doping profiles

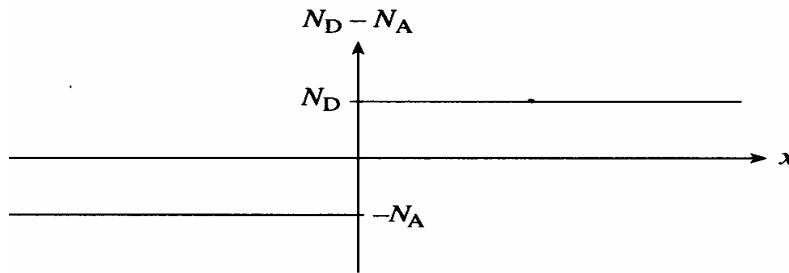
# Ideal step-junction doping profile



(a)

Figure 5.2

# Equilibrium energy band diagram for the pn junction

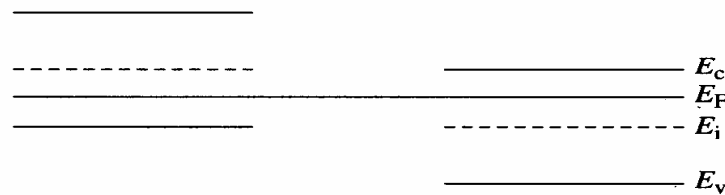


$$n = n_i \exp\left(\frac{E_F - E_i}{kT}\right)$$

$$p = n_i \exp\left(\frac{E_i - E_F}{kT}\right)$$

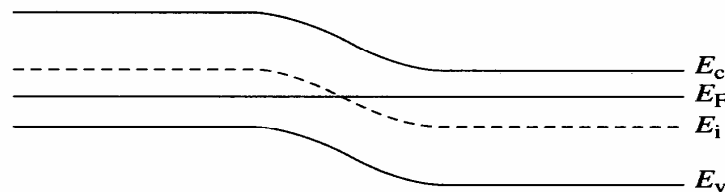


(a)



(b)

$E_F =$  same everywhere  
under equilibrium

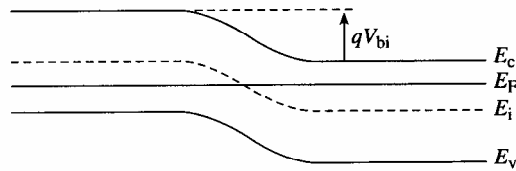


(c)

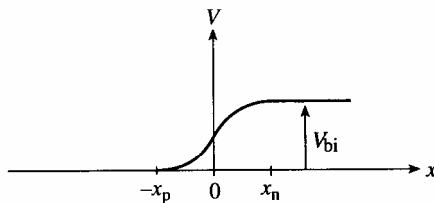
Figure 5.3

Join the two sides of the  
band by a smooth curve.

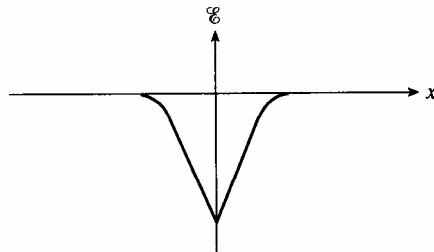
# Electrostatic variables for the equilibrium pn junction



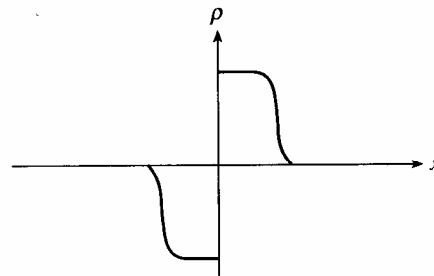
(a)



(b)



(c)



(d)

Figure 5.4

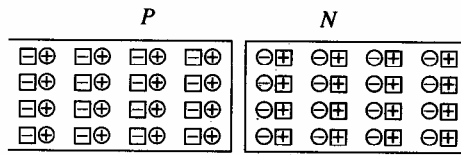
Potential,  $V = - (1/q) (E_C - E_{\text{ref}})$ . So, potential difference between the two sides (also called built-in voltage,  $V_{\text{bi}}$ ) is equal to  $-(1/q)(\Delta E_C)$ .

$$V = -\frac{1}{q} (E_C - E_{\text{ref}})$$

$$\mathcal{E} = \frac{1}{q} \frac{dE_C}{dx} = \frac{1}{q} \frac{dE_i}{dx}$$

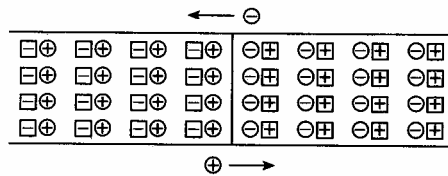
$$\left. \begin{aligned} \frac{d\mathcal{E}}{dx} &= \frac{\rho}{\epsilon} \end{aligned} \right\} \begin{aligned} \rho &= \text{charge density} \\ \epsilon &= K_s \epsilon_0 \end{aligned}$$

# Conceptual pn-junction formation

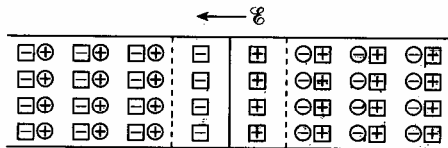


(a)

p and n type regions  
before junction formation

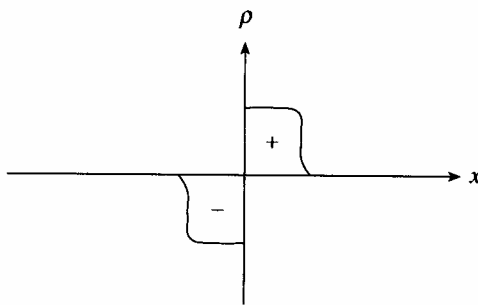


(b)



(c)

Holes and electrons will diffuse towards opposite directions, **uncovering ionized dopant atoms.** This will build up an electric field which will prevent further movement of carriers.



(d)

Figure 5.5

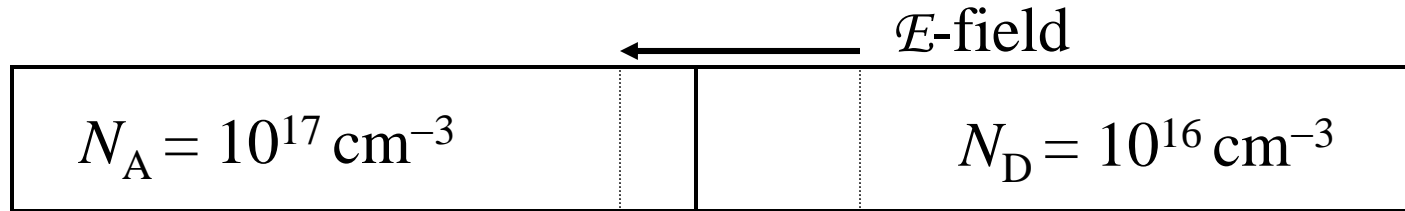
## The built-in potential, $V_{bi}$

When the junction is formed, electrons from the n-side and holes from the p-side will diffuse leaving behind charged dopant atoms. Remember that the dopant atoms cannot move! Electrons will leave behind positively charged donor atoms and holes will leave behind negatively charged acceptor atoms.

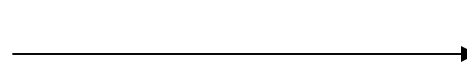
The net result is the build up of an electric field from the positively charged atoms to the negatively charged atoms, i.e., from the n-side to p-side. When steady state condition is reached after the formation of junction (how long this takes?) the net electric field (or the built in potential) will prevent further diffusion of electrons and holes. In other words, there will be **drift and diffusion currents such that net electron and hole currents will be zero.**

# Equilibrium conditions

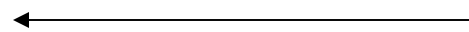
Under equilibrium conditions, the net electron current and hole current will be zero.



hole diffusion current

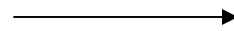


hole drift current

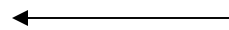


net = 0

electron diffusion current  
opposite to electron flux

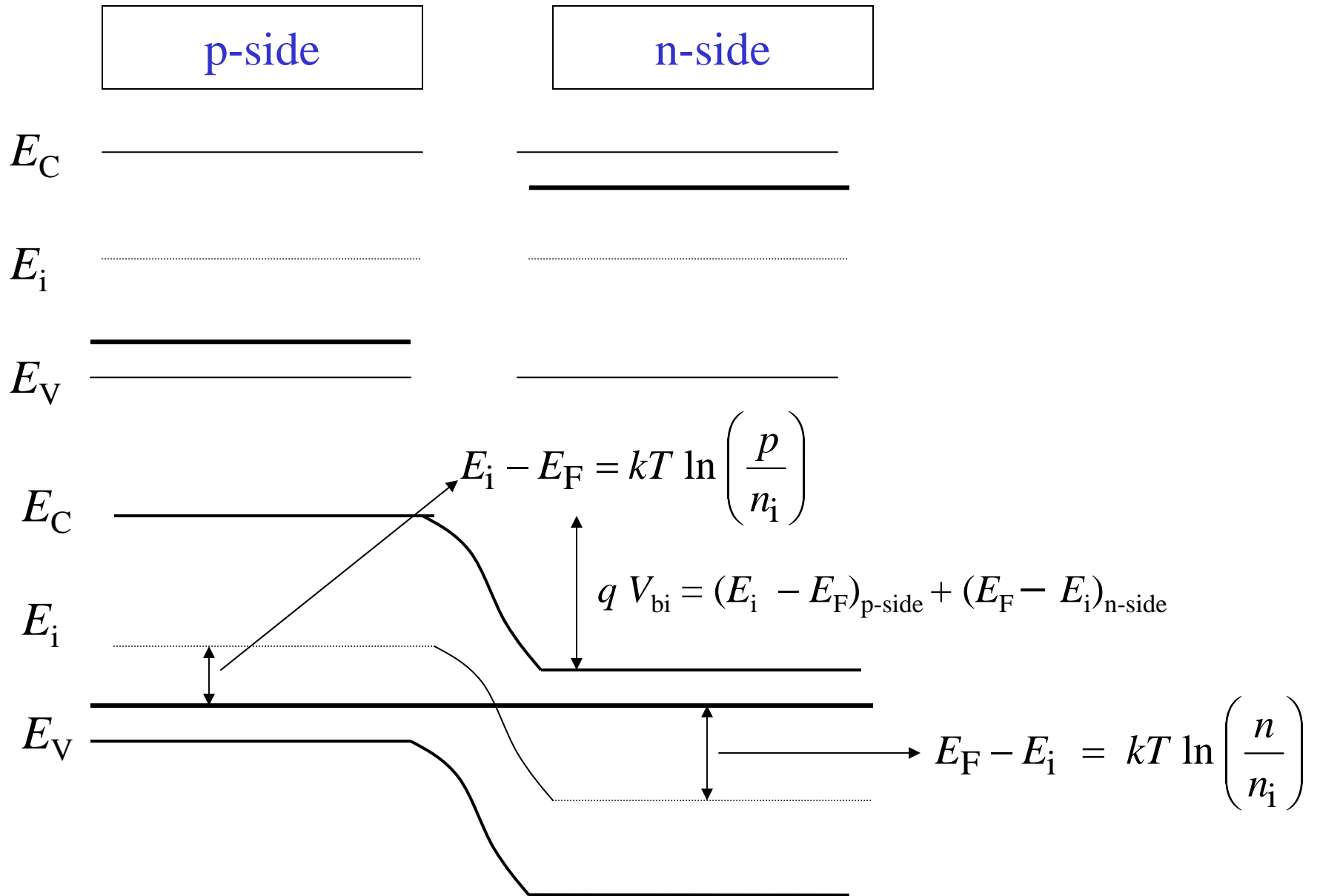


electron drift current  
opposite to electron flux



net = 0

# The built-in potential, $V_{bi}$



## The built-in potential, $V_{bi}$

The built-in potential,  $V_{bi}$ , measured in Volts, is numerically equal to the “shift” in the bands expressed in eV.

$$\begin{aligned} V_{bi} &= (1/q) \{ (E_i - E_F)_{p\text{-side}} + (E_F - E_i)_{n\text{-side}} \} \\ &= \frac{kT}{q} \ln \left( \frac{p}{n_i} \right) + \frac{kT}{q} \ln \left( \frac{n}{n_i} \right) \\ &= \frac{kT}{q} \ln \left( \frac{p_p n_n}{n_i^2} \right) \end{aligned}$$

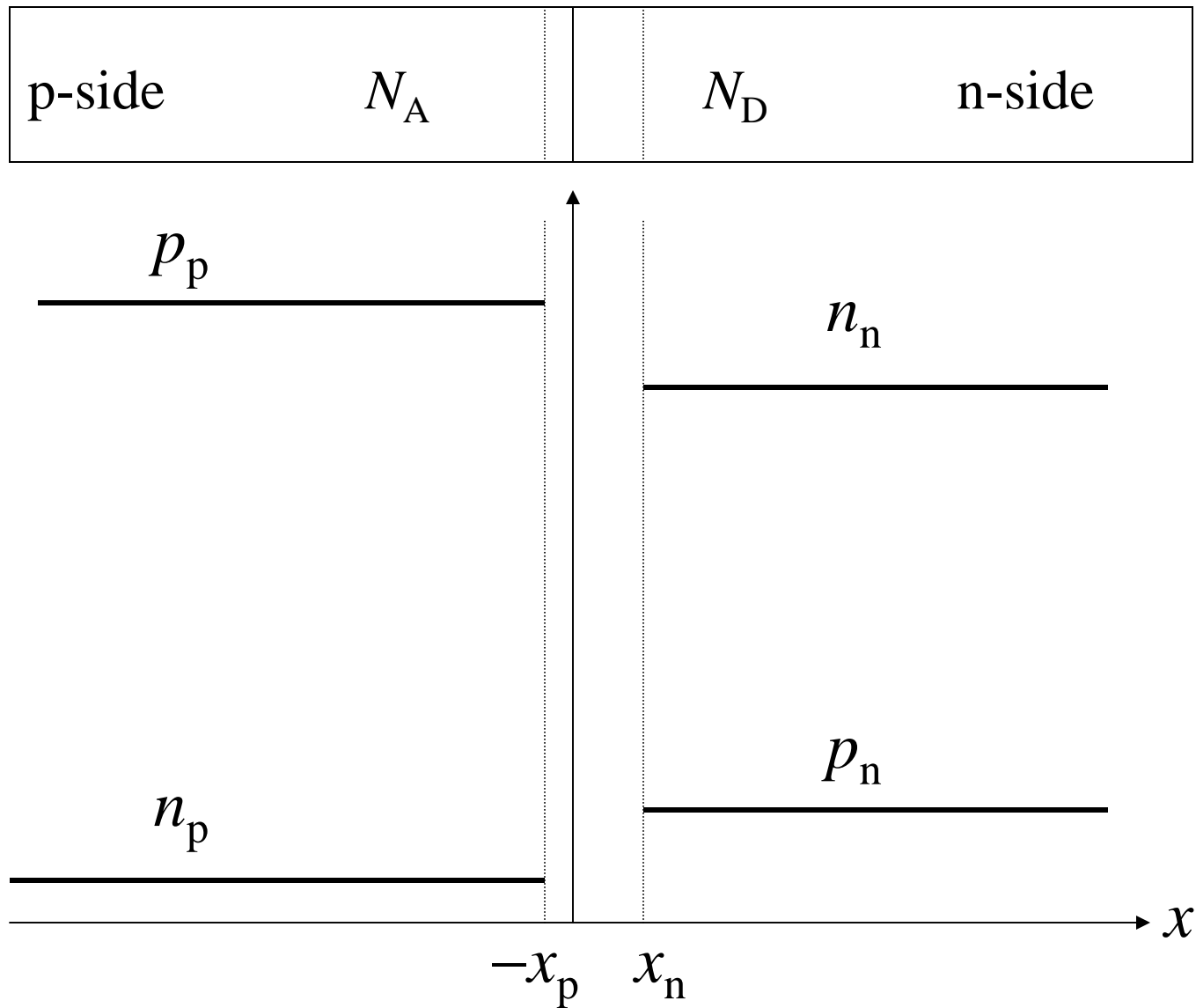
where  $p_p$  = hole – concentration on p – side

and  $n_n$  = electron – concentration on n – side

**An interesting fact:**

$$\frac{p_p}{p_n} = \frac{n_n}{n_p} = \exp \left( \frac{q V_{bi}}{kT} \right)$$

# Majority and minority carrier concentrations



# Built-in potential as a function of doping concentration for an abrupt p<sup>+</sup>n or n<sup>+</sup>p junction

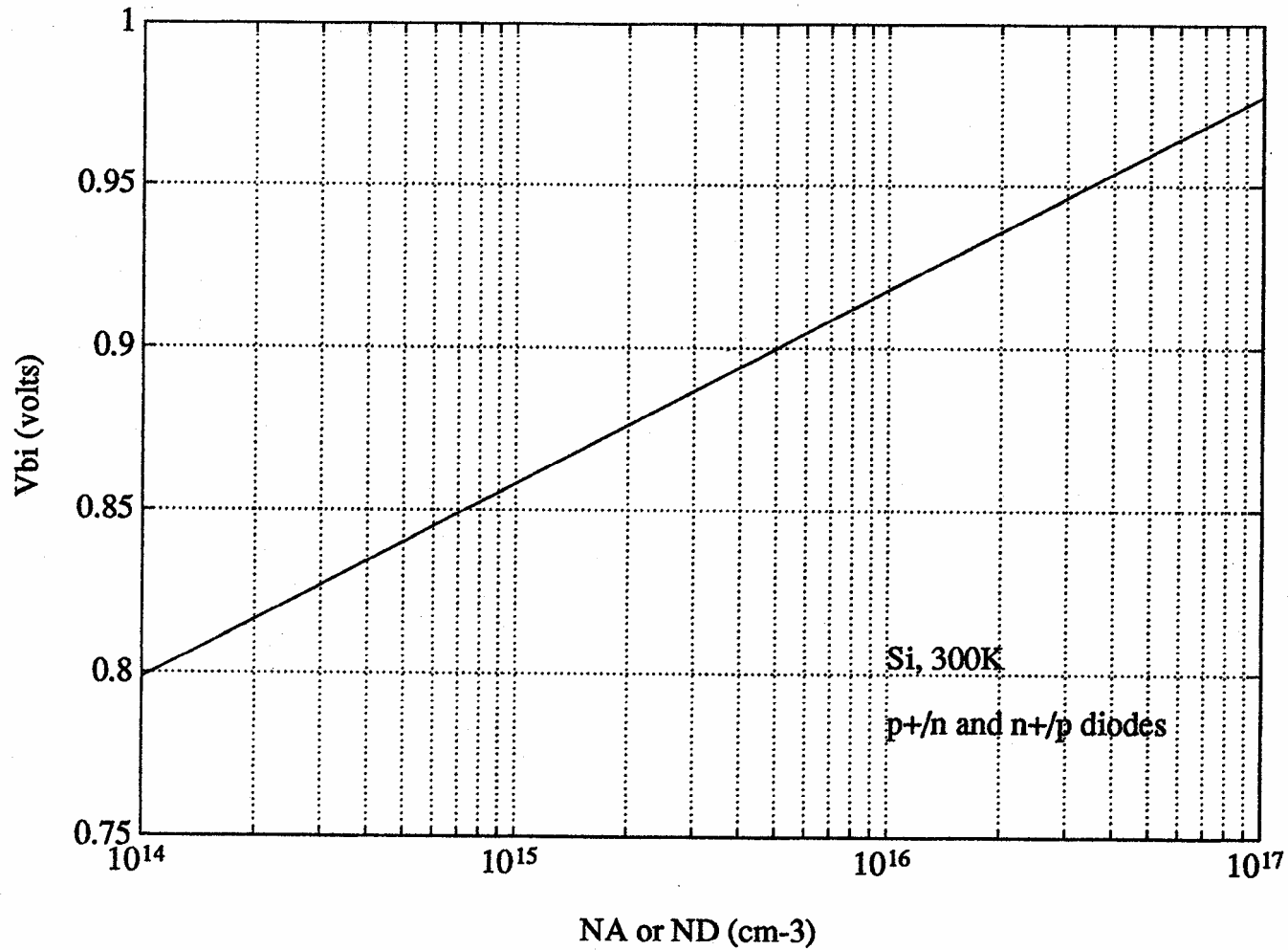
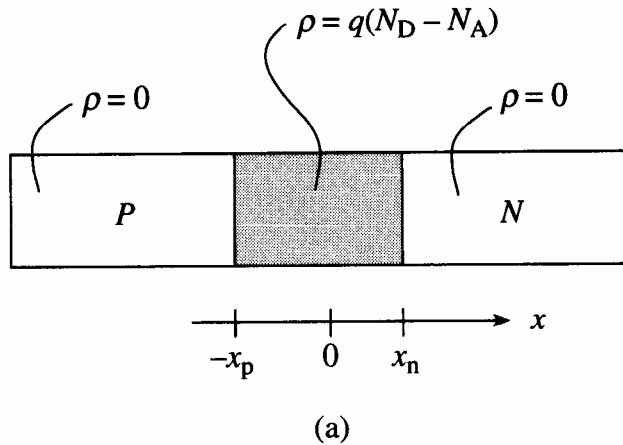


Figure E5.1

# Depletion approximation



$$\frac{dE}{dx} = \frac{\rho}{K_s \epsilon_0} \quad \text{Poisson equation}$$

$$= \begin{cases} \frac{q}{K_s \epsilon_0} (N_D - N_A) & \text{for } -x_p \leq x \leq x_n \\ 0 & \text{everywhere else} \end{cases}$$

*We assume that the free carrier concentration inside the depletion region is zero.*

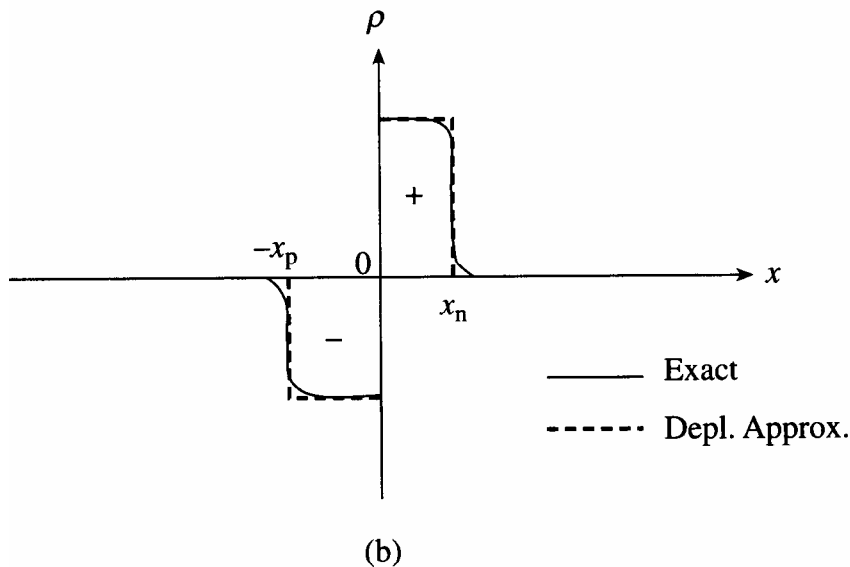


Figure 5.6

## Example 1

A p-n junction is formed in Si with the following parameters. Calculate the built-in voltage,  $V_{bi}$ .

$N_D = 10^{16} \text{ cm}^{-3}$	$N_A = 10^{17} \text{ cm}^{-3}$
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Calculate majority carrier concentration in n-side and p-side. Assume  $n_n = N_D = 10^{16} \text{ cm}^{-3}$  and  $p_p = N_A = 10^{17} \text{ cm}^{-3}$ .

$$V_{bi} = \frac{kT}{q} \ln \left( \frac{p_p n_n}{n_i^2} \right) = \frac{kT}{q} \ln \left( \frac{N_A N_D}{n_i^2} \right)$$

Plug in the numerical values to calculate  $V_{bi}$ .

## Example 2

A pn junction is formed in Si with the following parameters. Calculate the built-in voltage,  $V_{bi}$ .

$N_D = 2 \times 10^{16} \text{ cm}^{-3}$ $N_A = 10^{16} \text{ cm}^{-3}$	$N_A = 3 \times 10^{17} \text{ cm}^{-3}$ $N_D = 2 \times 10^{17} \text{ cm}^{-3}$
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Calculate majority carrier concentration in n-side and p-side.

$$n_n = \text{“effective } N_D \text{”} = 10^{16} \text{ cm}^{-3}. \quad p_p = \text{“effective } N_A \text{”} = 10^{17} \text{ cm}^{-3}.$$

$$V_{bi} = \frac{kT}{q} \ln \left( \frac{p_p n_n}{n_i^2} \right) = \frac{kT}{q} \ln \left( \frac{N_A N_D}{n_i^2} \right) \quad \left. \vphantom{V_{bi}} \right\} \text{ Here } N_A \text{ and } N_D \text{ are “effective” or net values.}$$

Plug in the numerical values to calculate  $V_{bi}$ .